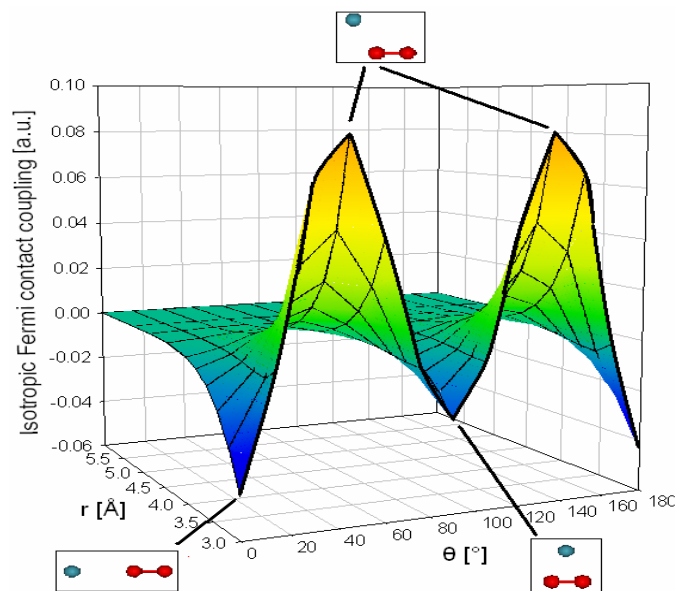


# Intermolecular hyperfine tensor for $\text{Xe-O}_2$ and density and temperature dependence of Xe chemical shifts in oxygen gas



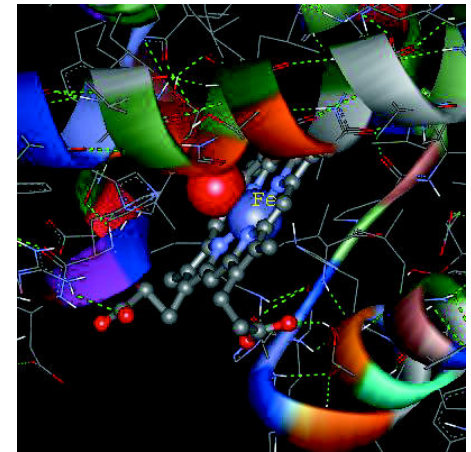
**Lela Vukovic and  
Cynthia J. Jameson**  
ENC 2005



# Paramagnetic centers play important roles

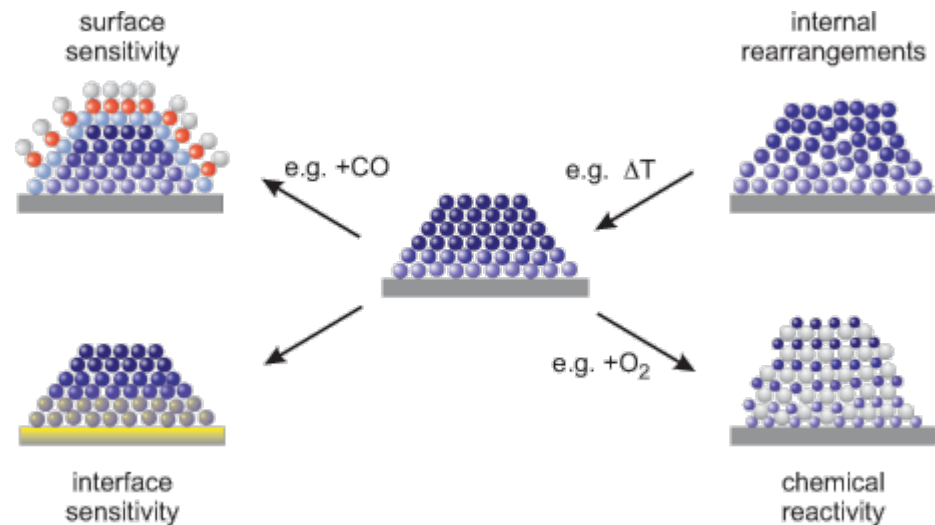
## haem proteins:

- catalysis – peroxidases, catalases
- electron transfer – cytochromes
- oxygen transport and storage – globins
- nitric oxide transport – nitrophorin



[www.linux-magazin.de](http://www.linux-magazin.de)

Structure and particularly the interfacial characteristics of the particles are strongly correlated with their magnetic properties.

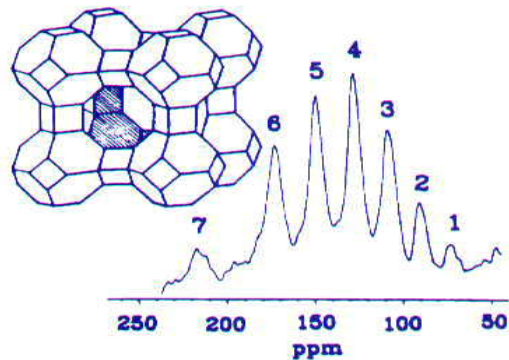


Paramagnetic centers at solid surfaces

(from Chemical Physics Magnetic Resonance Group, Fritz-Haber Institute)

# Xe is a good probe of porous materials

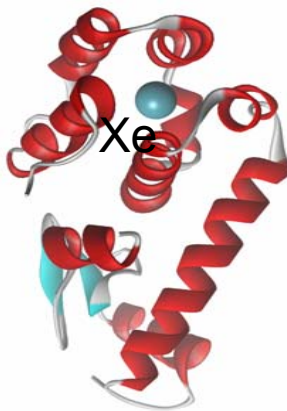
distributions of adsorbed atoms among cages



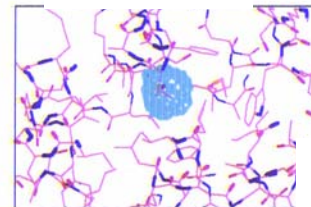
C. J. Jameson, A. K. Jameson, R. Gerald, A. C. de Dios, J. Chem. Phys. 97, 417 (1992).

Xe in zeolite NaA

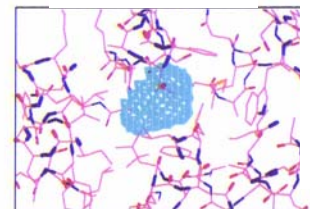
small differences in amino acids near the pocket



Wild Type

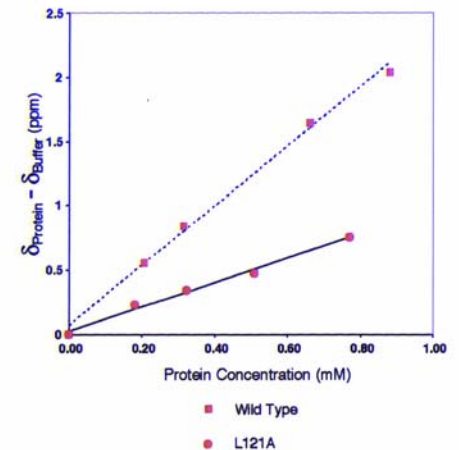


L121A



$\delta_{\text{Protein}} - \delta_{\text{Buffer}}$  (ppm)

Xe NMR Data



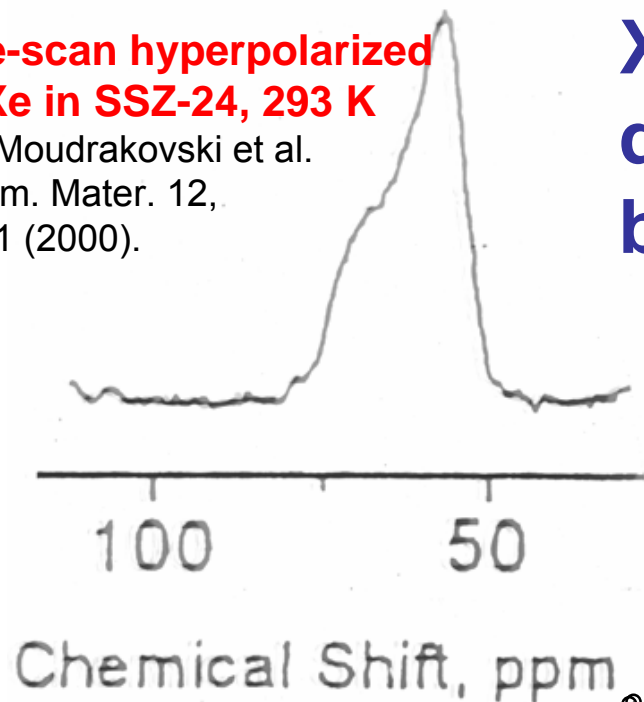
S.M. Rubin, M.M. Spence, B.M. Goodson, D.E. Wemmer, A. Pines.  
Proc. Natl. Acad. Sci. **97** (17) 9472-9475 (2000)

Xe in protein pocket

## EXPERIMENT

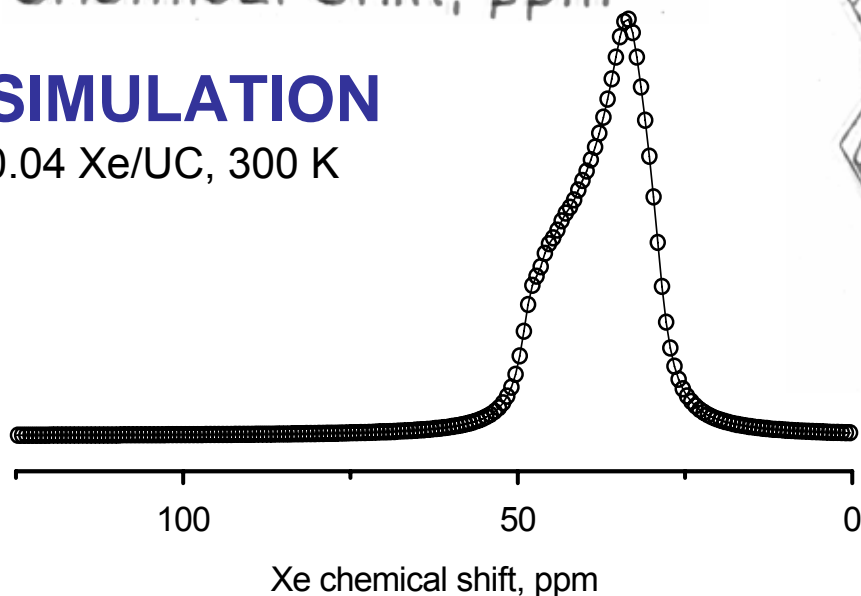
one-scan hyperpolarized  
 $^{129}\text{Xe}$  in SSZ-24, 293 K

I.L. Moudrakovski et al.  
Chem. Mater. 12,  
1181 (2000).

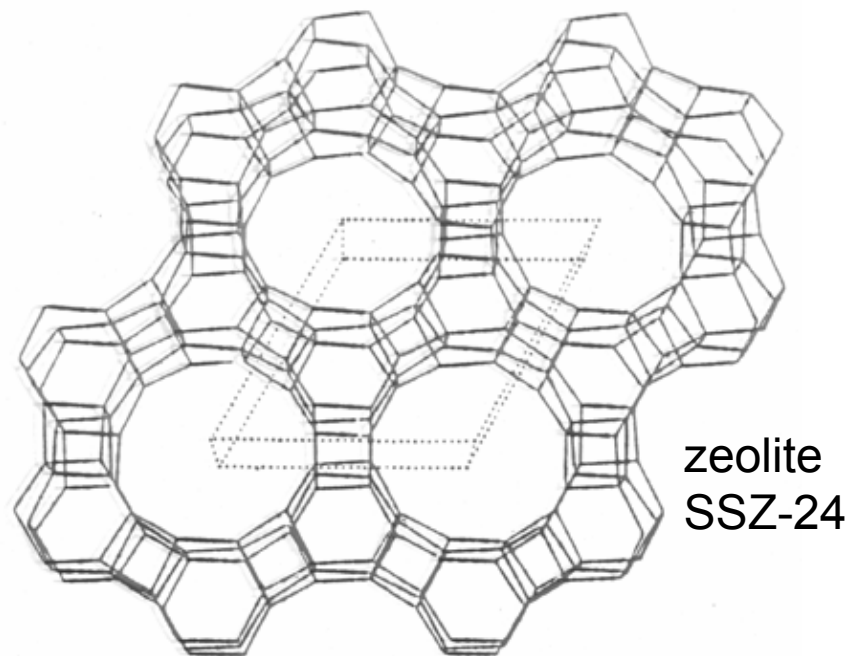


## SIMULATION

0.04 Xe/UC, 300 K



Quantitative understanding of  
Xe chemical shifts in porous  
diamagnetic materials has  
been demonstrated

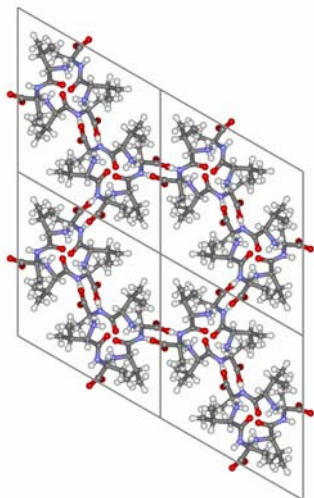


GCMC simulation uses the same  
shielding and potential functions as  
for Xe in silicalite.

C.J. Jameson, JACS, 126, 10450 (2004)

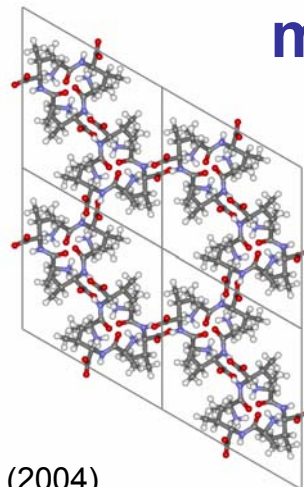
in channels of

molecular crystals



L-Val-L-Ala

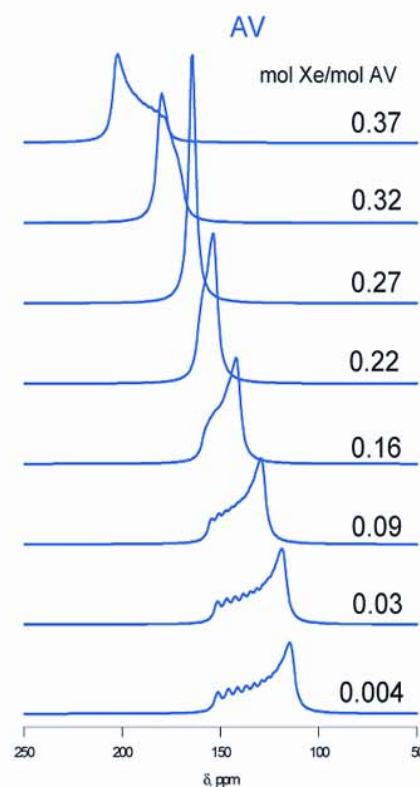
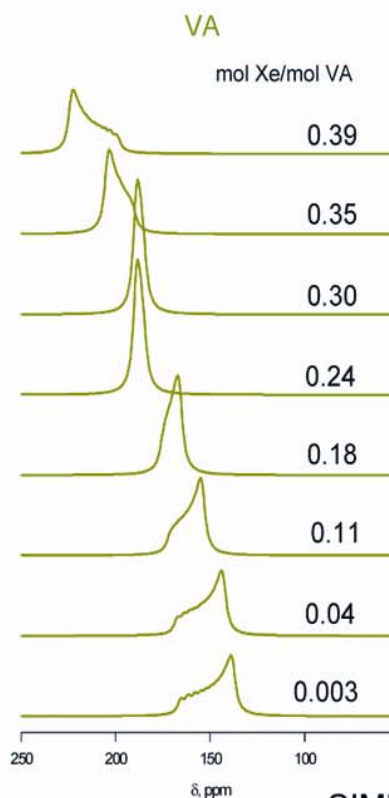
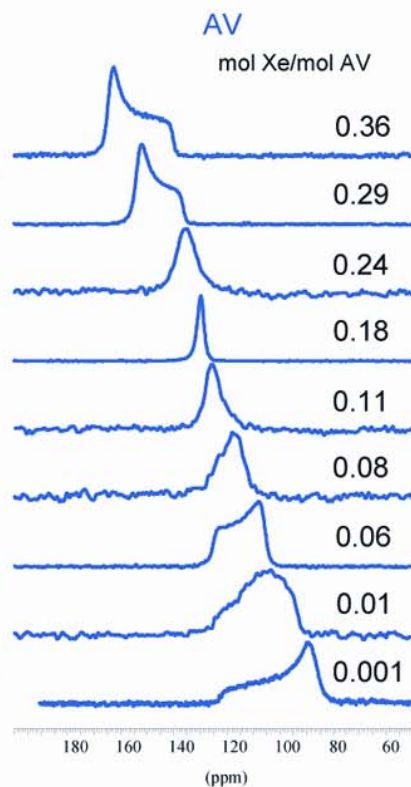
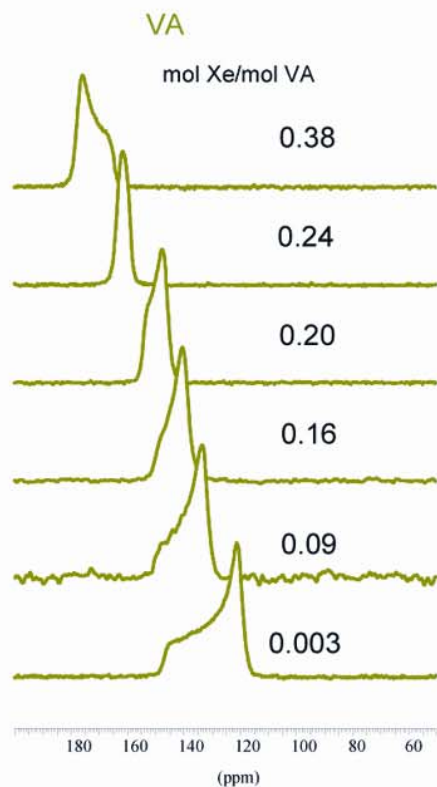
VA



L-Ala-L-Val

AV

PNAS 101, 17924 (2004)

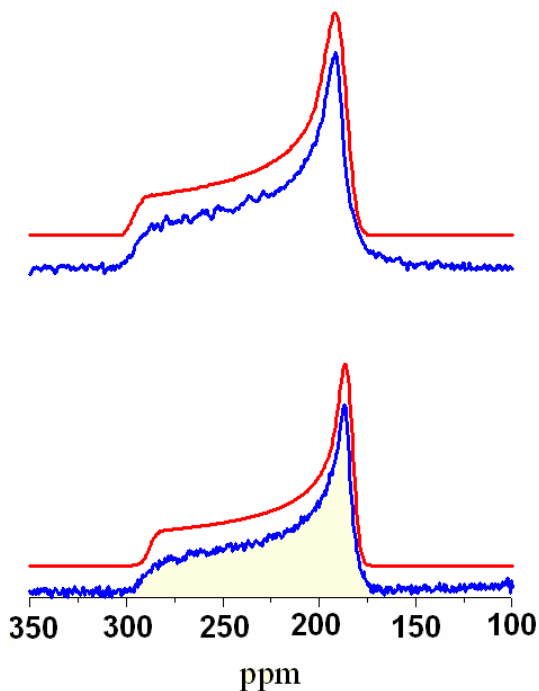


EXPERIMENT

SIMULATION

Xe NMR spectra

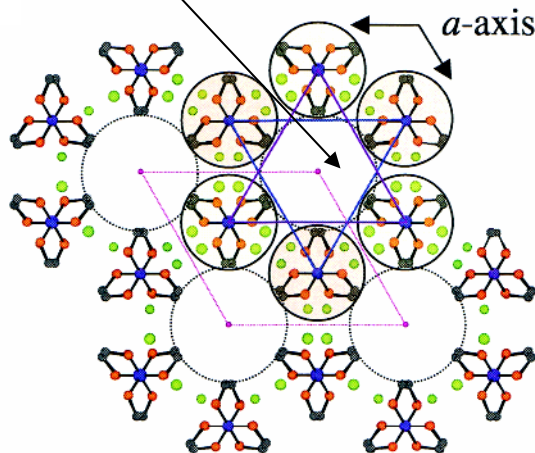
# Xe as a probe of porous materials with paramagnetic centers?



Wasylishen et al., 2004

$[\text{Rh}(\text{en})_3]\text{Cl}_3$  crystal  
 $[\text{Co}(\text{en})_3]\text{Cl}_3$  crystal  
 diamagnetic

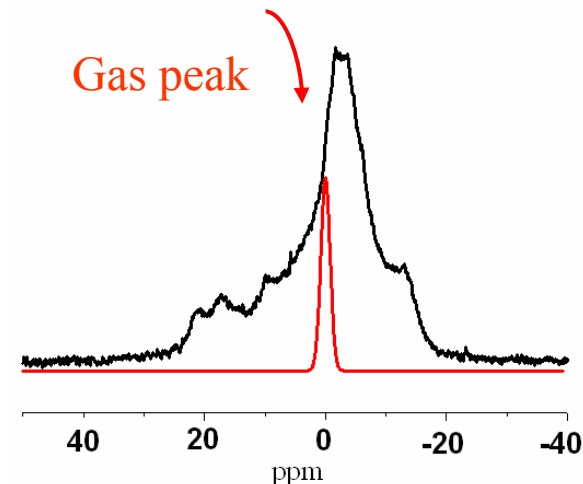
Xe inside these channels



same crystal structure

Ueda et al.  
 J.Phys. Chem. B 107, 180 (2003)

1000 scans using ~10%  
 hyperpolarized xenon gas



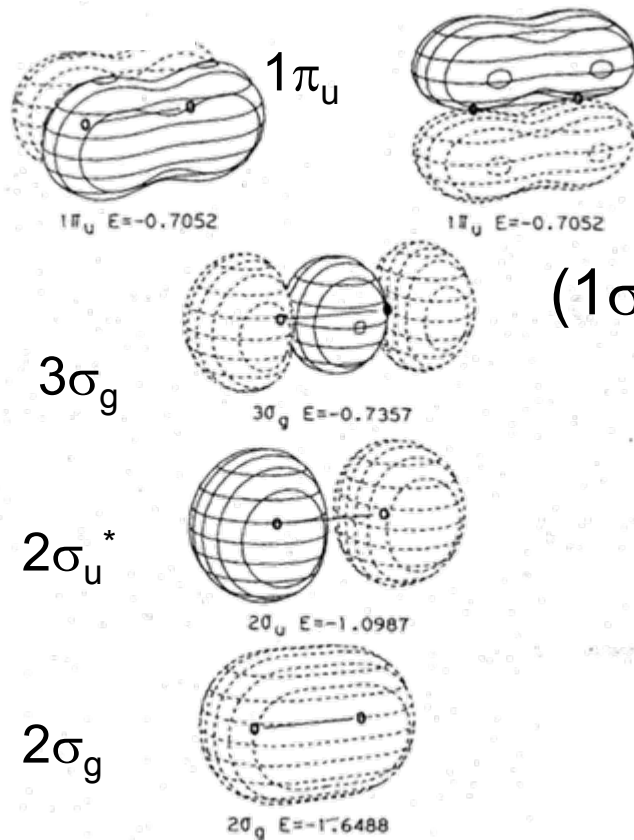
Wasylishen et al., 2004

$[\text{Cr}(\text{en})_3]\text{Cl}_3$  crystal  
 paramagnetic

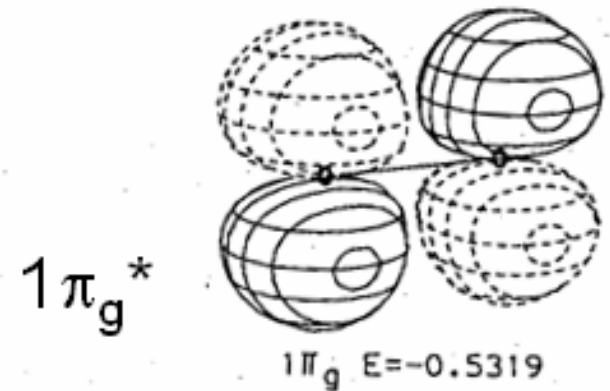
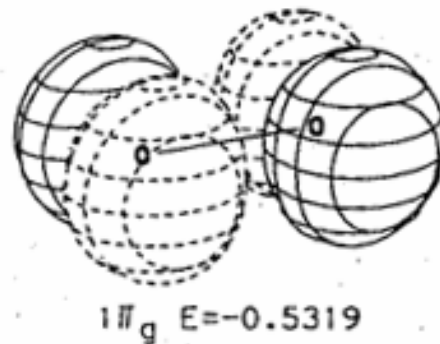
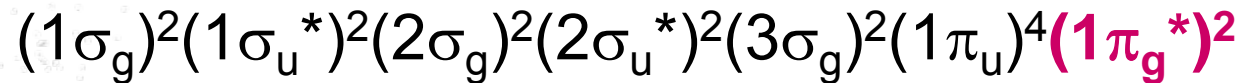
Xe can tell the difference!

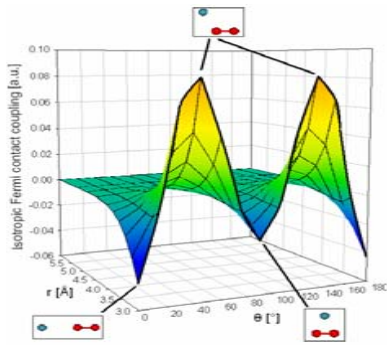
# Study Xe in the presence of a paramagnetic center

## Our model system:



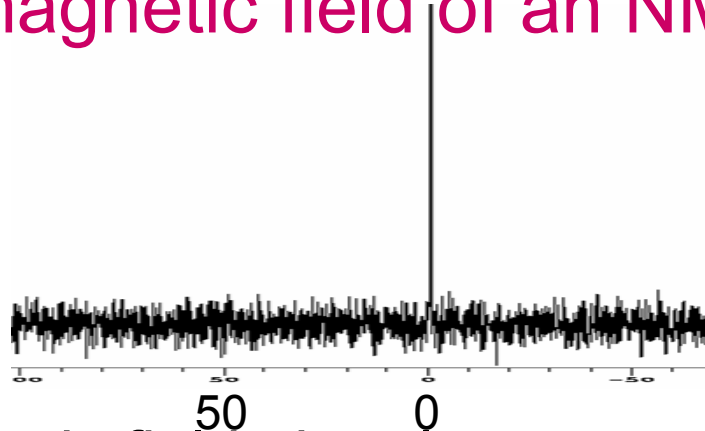
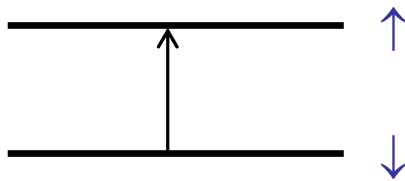
$O_2$  molecule has two unpaired electrons



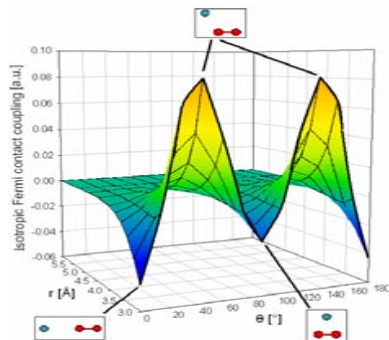


# The Xe atom in an NMR experiment

The Xe nucleus is a magnetic dipole that has two states in the presence of the magnetic field of an NMR spectrometer.

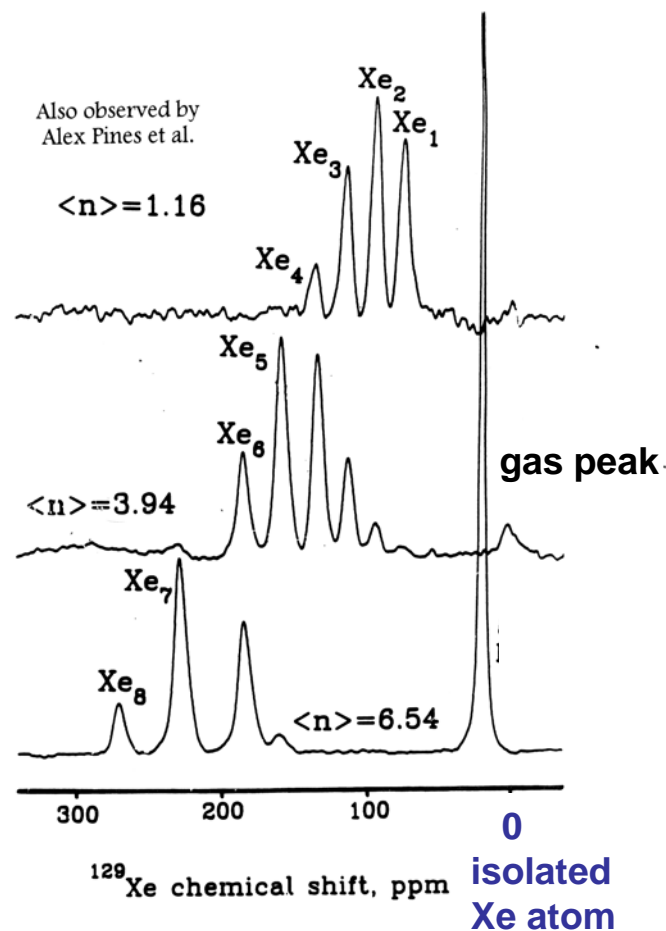


In the presence of a magnetic field, the electrons of Xe atom provide a diamagnetic response, that is, a shielding effect, the characteristic “**NMR chemical shift**” of the free Xe atom. A sample of very very low density Xe gas serves as an NMR sample of individual non-interacting Xe atoms, a reference which we can call 0 ppm on the Xe chemical shift scale.



# The Xe atom in an NMR experiment

The Xe chemical shift is very sensitive to the presence of neighbors.



# Xe atom in a magnetic field in the presence of a diamagnetic molecule, N<sub>2</sub>

A Xe atom in the presence of a diamagnetic molecule such as N<sub>2</sub> experiences an intermolecular deshielding effect that can be calculated quantum mechanically.

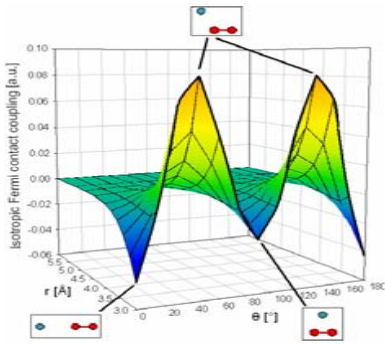
This intermolecular response leads to a positive Xe chemical shift compared to the reference Xe atom. The chemical shift increases with density, with increasing probability of encounters between Xe and N<sub>2</sub> in the gas.

$$\sigma(T, \rho_{N_2}) = \sigma(\text{Xe atom}) + \sigma_1(T) \rho_{N_2} + \dots$$

The **density coefficient**  $\sigma_1(T)$  is the ‘per molecule average’ over such probabilities, if  $\rho_{N_2}$  is a number density:

$$\sigma_1(T) = 2\pi \int \{\sigma(R, \theta) - \sigma(\infty)\} e^{-V(R, \theta)/kT} R^2 dR$$

# Xe in the presence of a paramagnetic molecule $O_2$

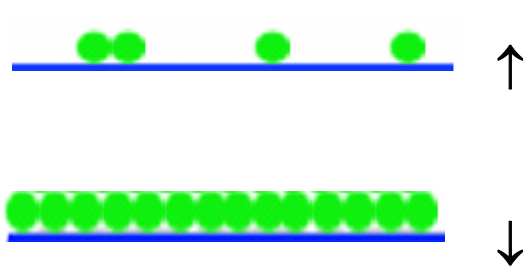


There is an additional interaction for Xe in  $O_2$  because of the unpaired electrons:

**the hyperfine shift**

# The mechanism of the hyperfine shift

The average electron spin magnetic moment (along the  $B_0$  direction) of unpaired spins depends on  $B_0$  and the temperature



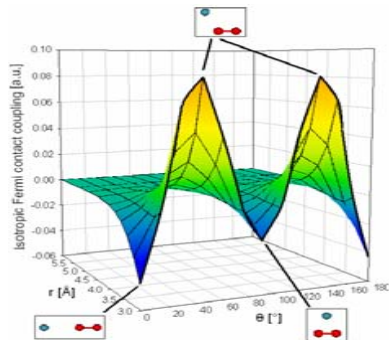
The **relative populations** of electron spin states at a given temperature determine the **average** electron spin magnetic moment along  $B_0$

The interaction between the Xe nucleus and the unpaired electrons of  $O_2$  molecule is:

$$\text{Interaction Energy} = \mathcal{A} \mathbf{S} \cdot \mathbf{I}$$

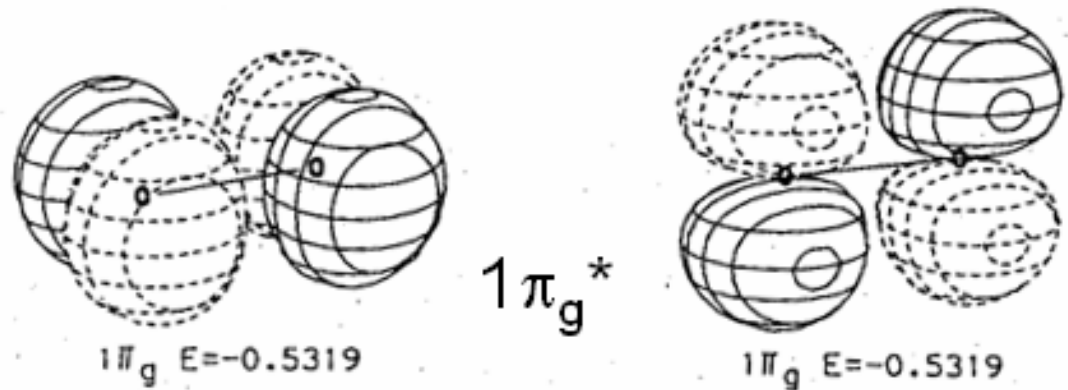
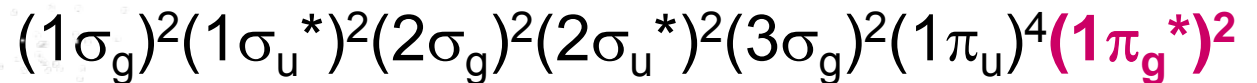
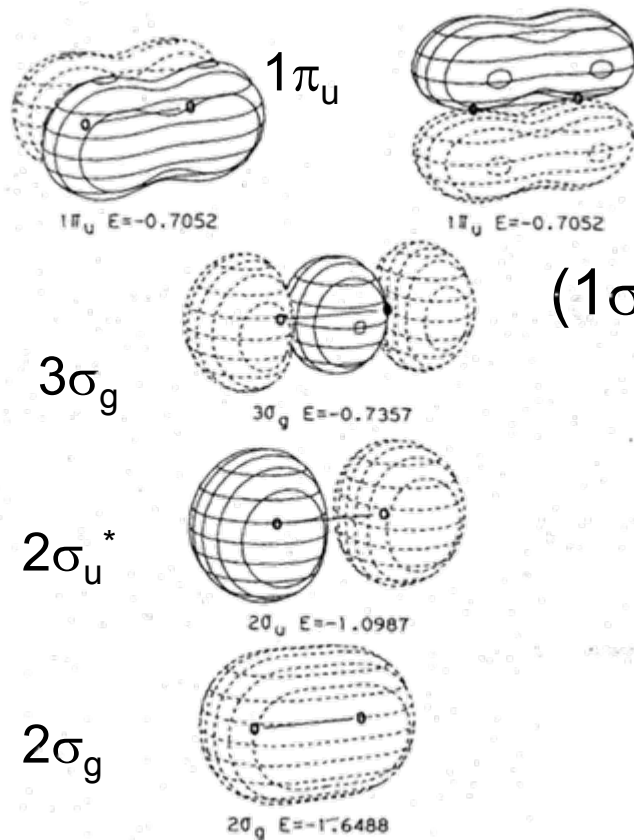
providing a fractional correction to the local magnetic field at the nucleus:

$$\Delta B/B_0 = -(g_e \mu_B / 3kT) \cdot (\mathcal{A} / g_{Xe} \mu_N)$$



# O<sub>2</sub>, a paramagnetic molecule

O<sub>2</sub> molecule has two unpaired electrons



# Quantum Mechanical Calculations

For the system  $\text{Xe}@O_2$ , we calculate quantum mechanically

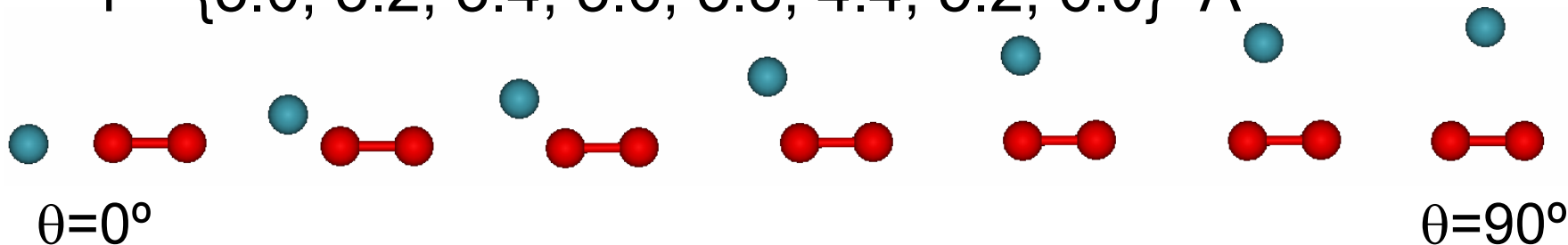
the net electron spin density at the position of the Xe nucleus, or the  $\mathcal{A}$  tensor itself.

We do this at various distances  $R$  and angles  $\theta$ .

## Configurations:

$$\theta = \{0, 15, 30, 45, 60, 75, 90\}^\circ$$

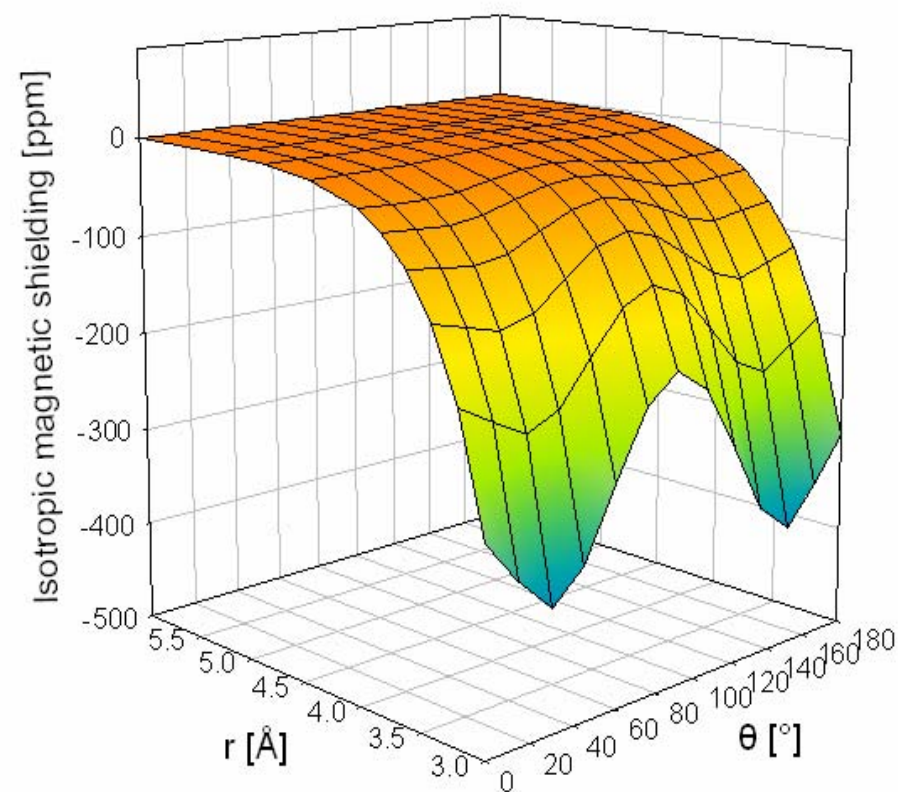
$$r = \{3.0, 3.2, 3.4, 3.6, 3.8, 4.4, 5.2, 6.0\} \text{ \AA}$$



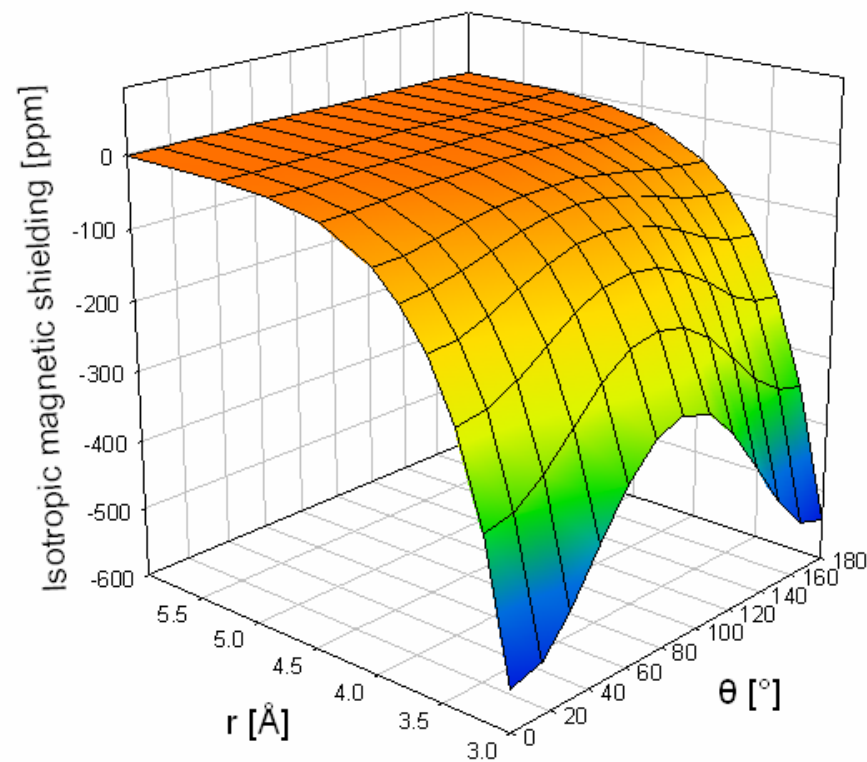
Using 240 basis functions on Xe and 80 on  $O_2$ .

Quantum mechanical calculations of **intermolecular shielding**  
Xe atom in the presence of O<sub>2</sub> experiences  
**a very similar intermolecular deshielding effect**  
as for Xe in diamagnetic N<sub>2</sub>

- Distance and angle dependence
- Magnitudes

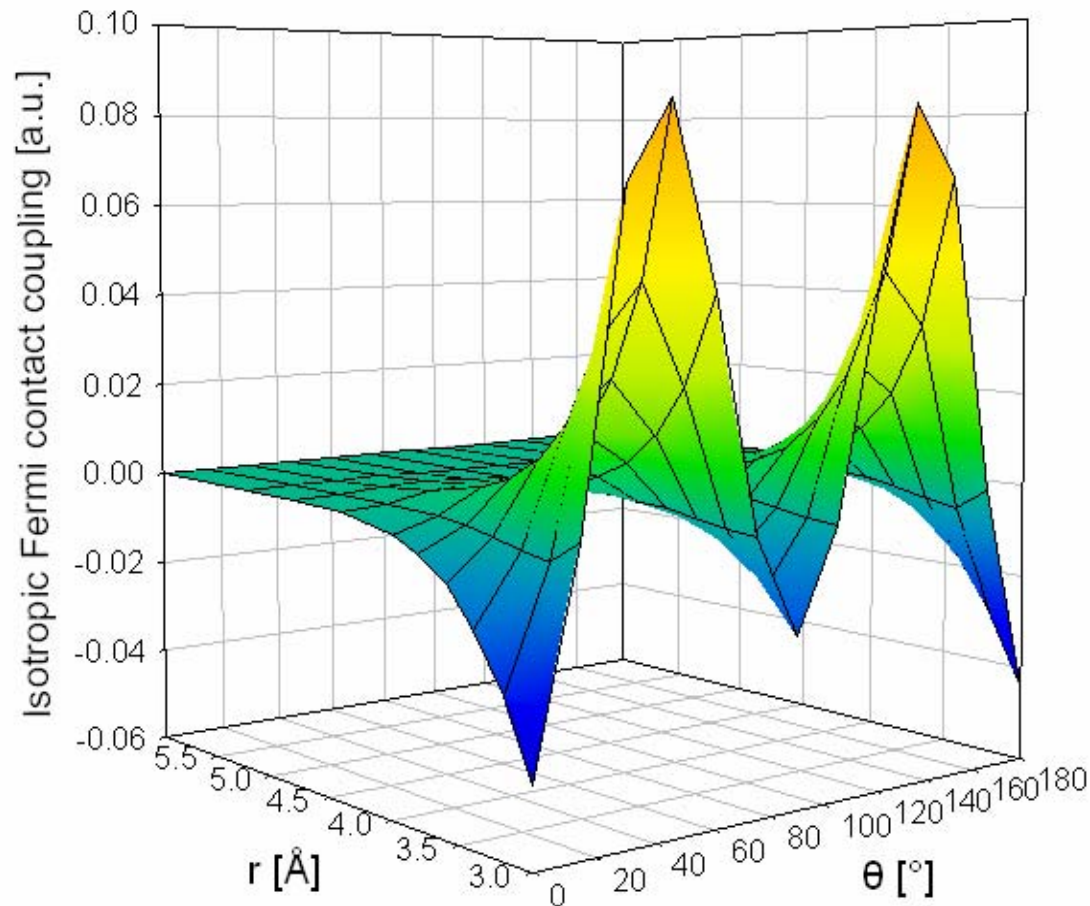


Xe@O<sub>2</sub>



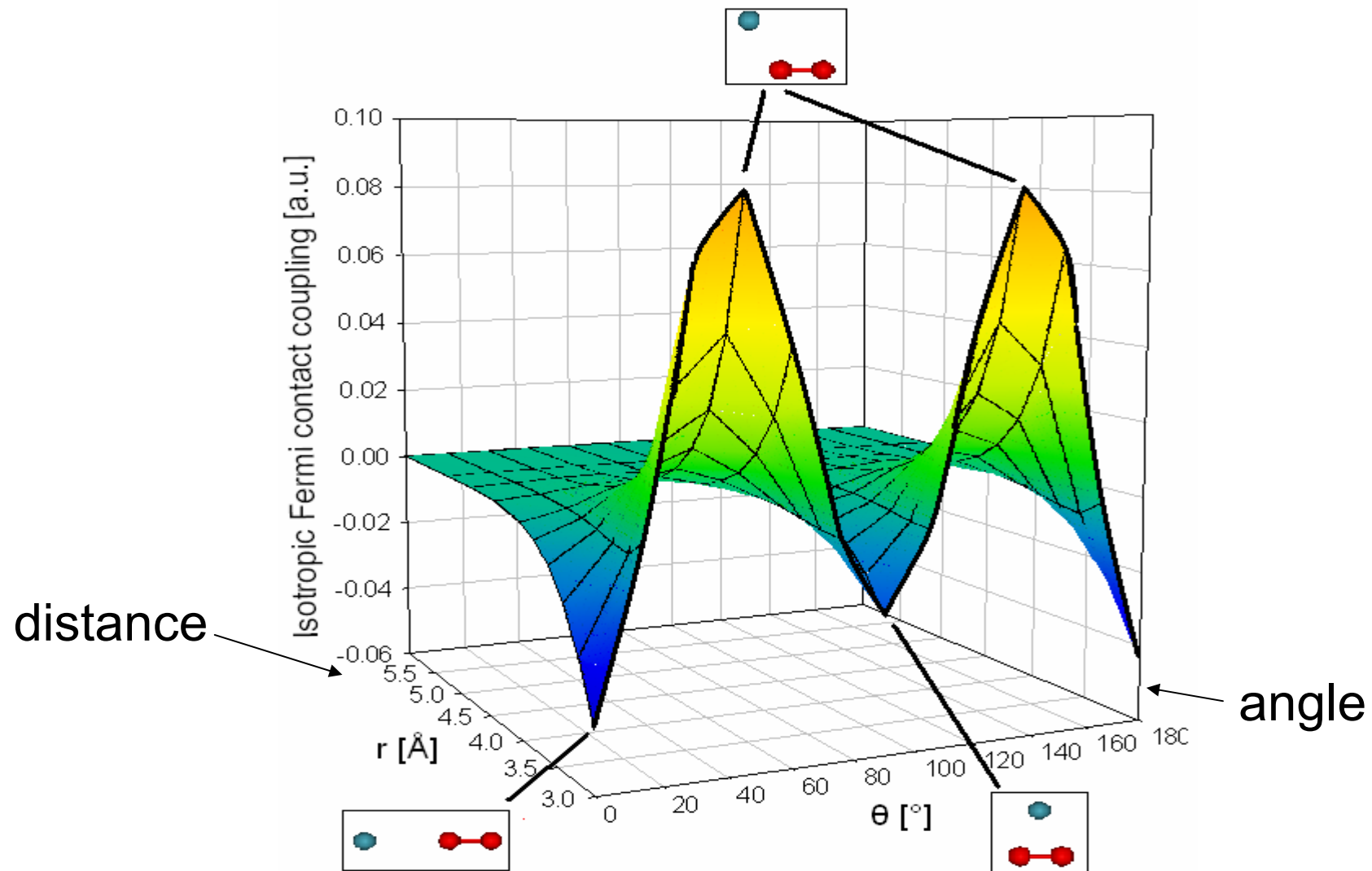
Xe@N<sub>2</sub>

# Quantum mechanical calculations of the **hyperfine interaction** for Xe@O<sub>2</sub>



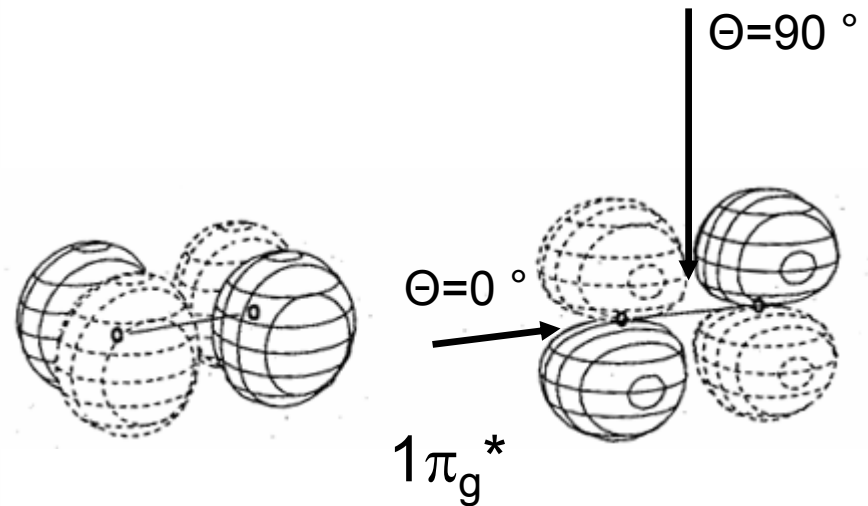
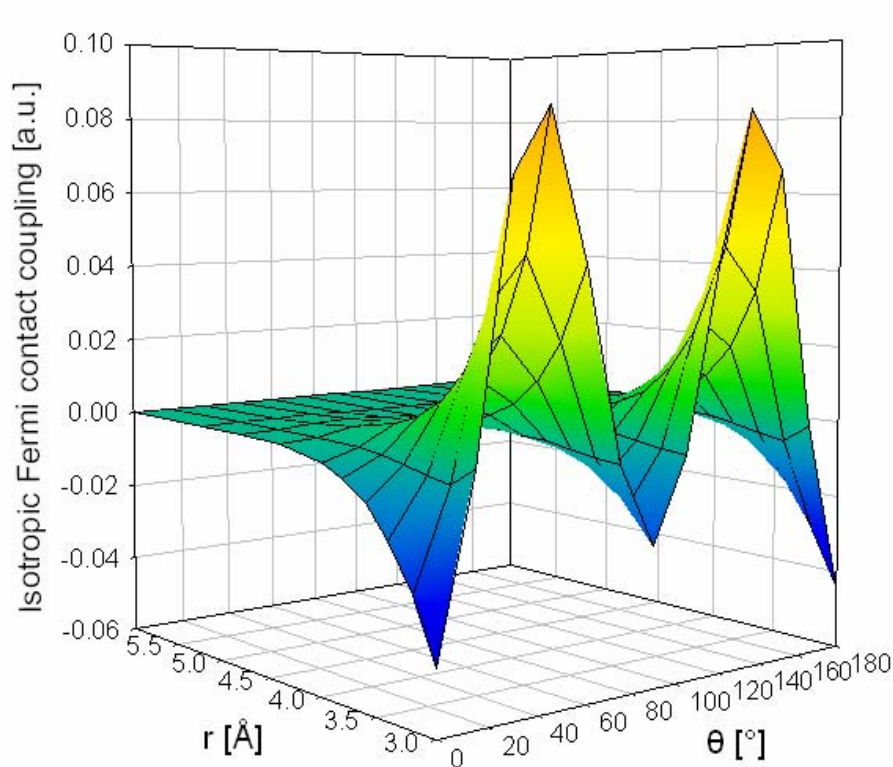
**Isotropic Fermi contact coupling surface**

Xe can tell **how far away** is the paramagnetic center and **how it is oriented** relative to the Xe position in the magnetic field!



Isotropic part of the hyperfine tensor

# Isotropic part of the hyperfine tensor for Xe@O<sub>2</sub>



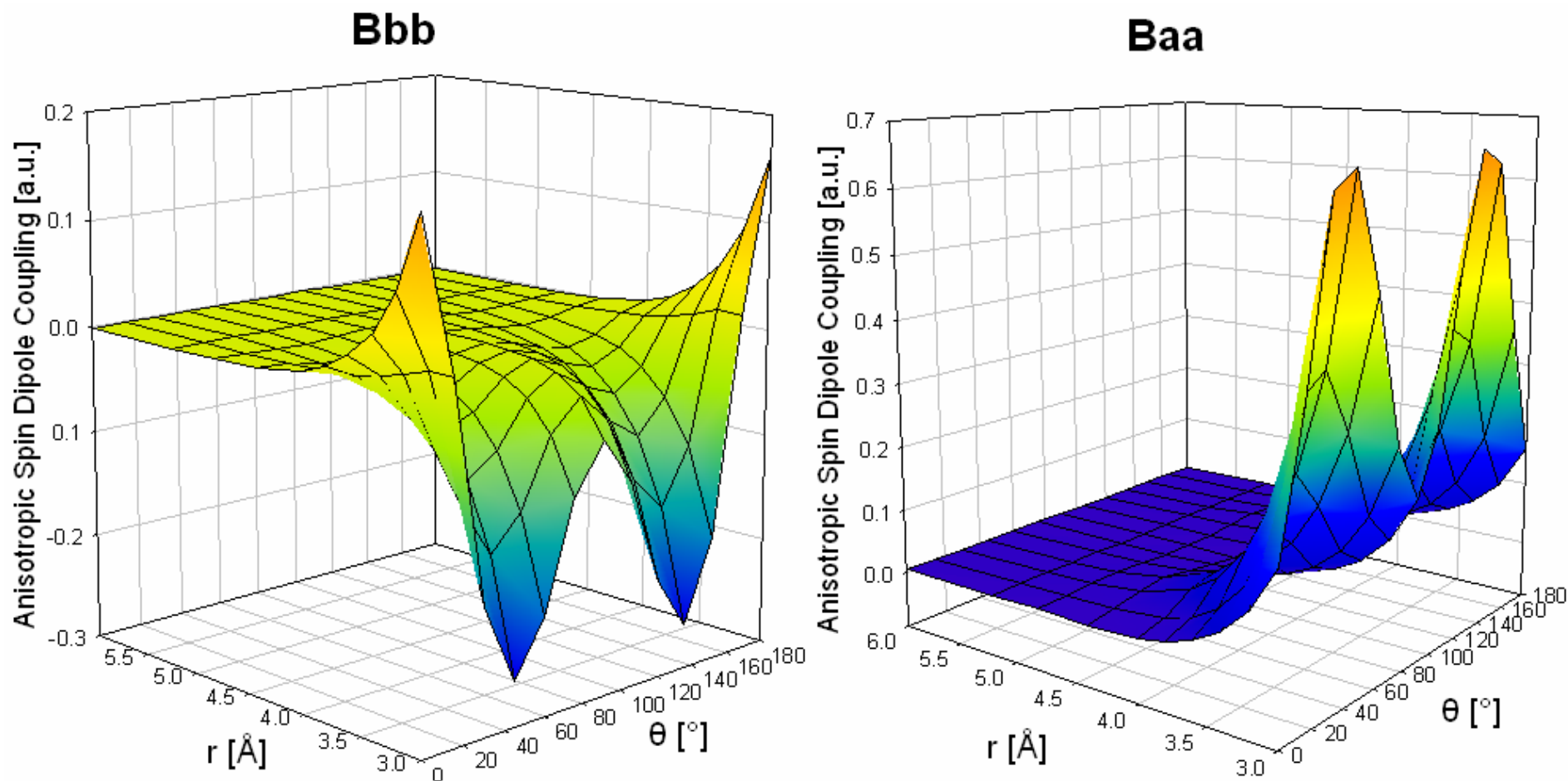
Angle dependence:

maximum at  $\theta = 45^\circ$

minimum at  $\theta = 0^\circ, 90^\circ$

Minima occur at the angles which coincide with the nodal planes of the O<sub>2</sub> molecule.

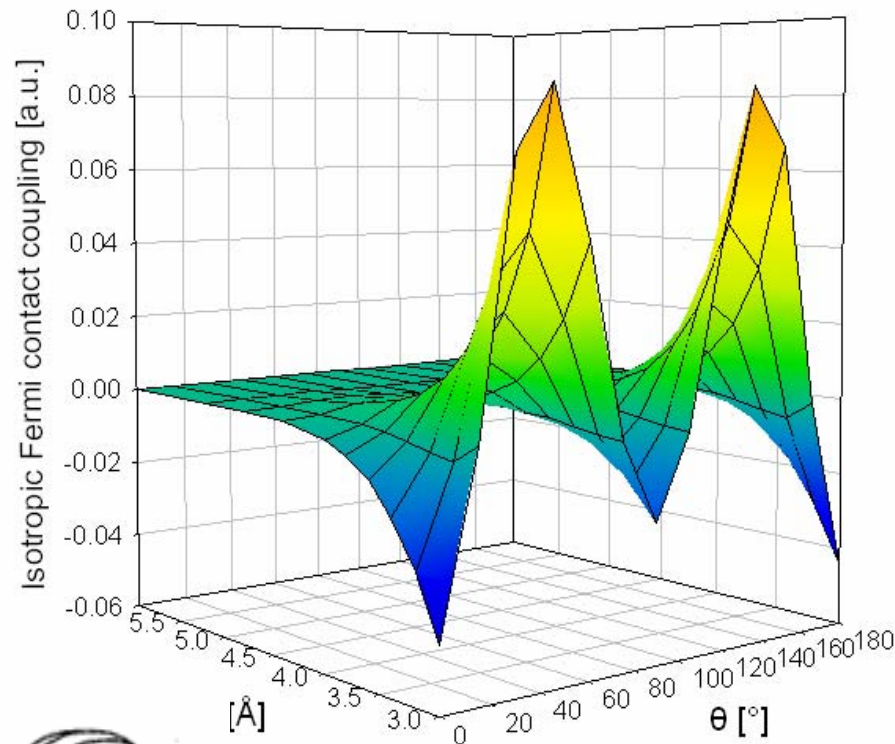
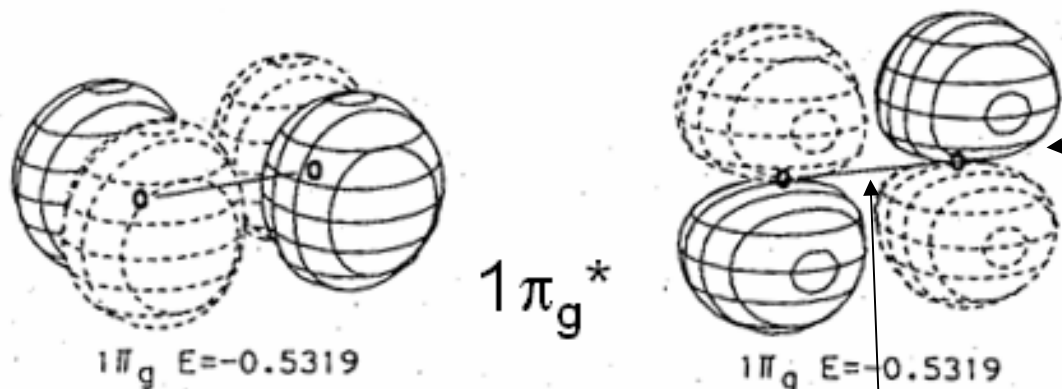
# The dipolar part of the hyperfine tensor:



Two of the principal components of the traceless tensor. **The dipolar part is relevant to line shape of Xe in channels with paramagnetic centers.**

# Why the angle dependence?

Simple picture: Unpaired electron spins reside, one apiece, in the  $1\pi_g^*$  molecular orbital of  $O_2$ :



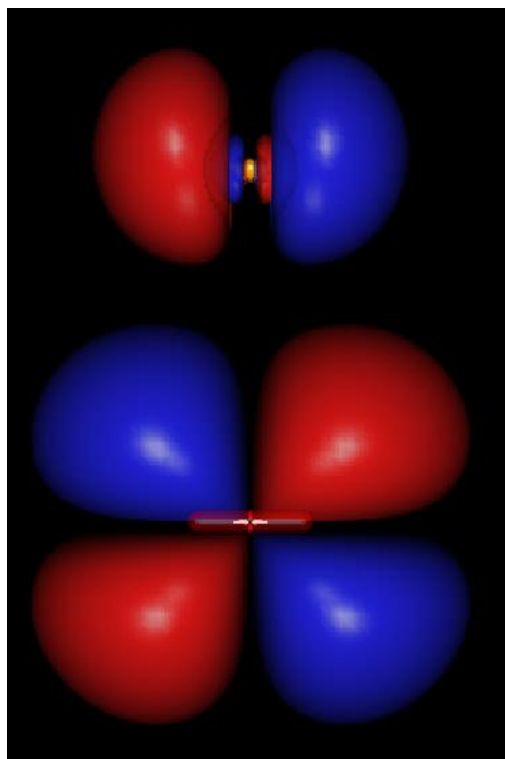
Xe ( $\theta=0^\circ$ )

At  $0^\circ$  and  $90^\circ$  the Xe encounters nodes of this  $O_2$  molecular orbital where the spin density is nil

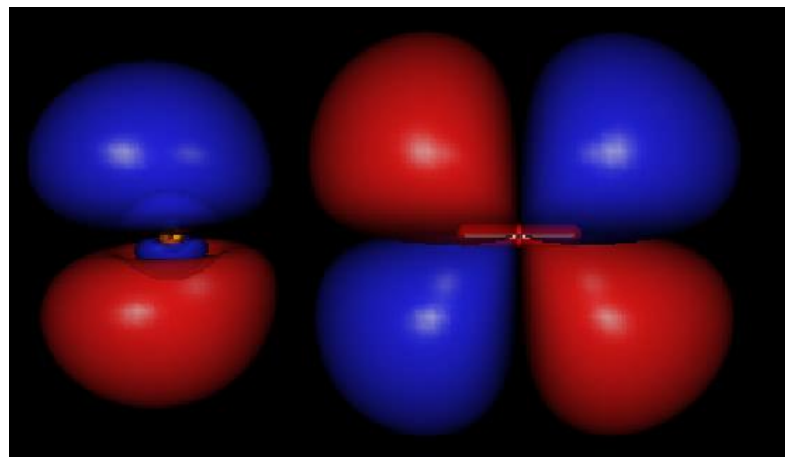
**BEST at  $45^\circ$  !!**

Xe ( $\theta=90^\circ$ )

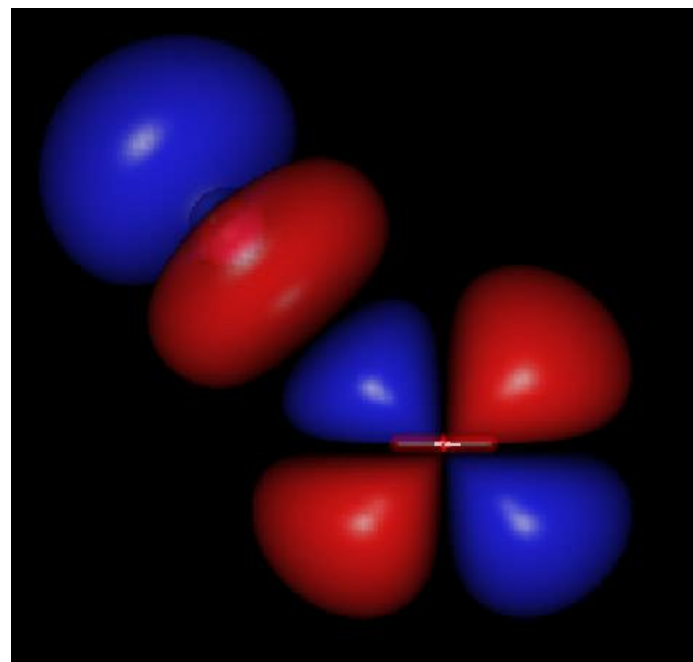
Largest contribution to spin density comes from this MO in the  $\text{XeO}_2$  complex.



$\text{Xe-O}_2$ , 90° orientation,  $r_{\text{Xe-O}_2} = 4.0 \text{ \AA}$

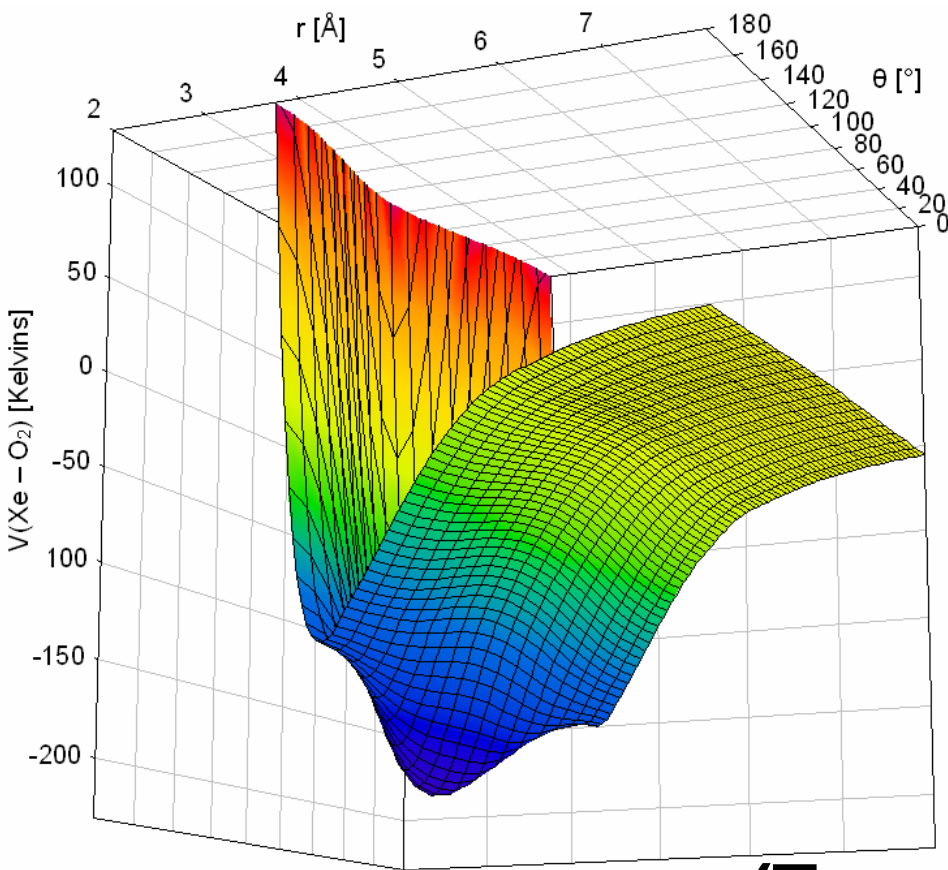


$\text{Xe-O}_2$ , 180° orientation,  $r_{\text{Xe-O}_2} = 4.0 \text{ \AA}$



$\text{Xe-O}_2$ , 45° orientation,  $r_{\text{Xe-O}_2} = 4.0 \text{ \AA}$

# Testing the calculations



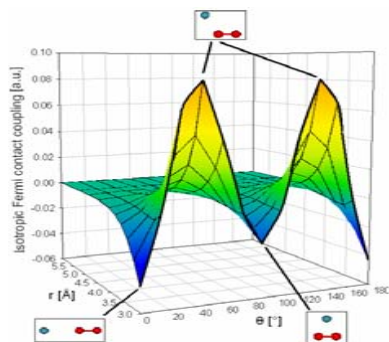
High quality potential function from:  
V. Aquilanti, D. Ascenzi, D. Cappelletti, M. de Castro,  
and F. Pirani, J. Chem. Phys. 109, 3898 (1998)

Using this potential energy function,  $V(\mathbf{R}, \theta)$ , the hyperfine tensor and shielding values are averaged over all configurations. The average is expressed as the density coefficient of the chemical shift.

$$\sigma(\mathbf{T}, \rho_{\text{O}_2}) - \sigma(\text{Xe atom}) = \sigma_1(\mathbf{T}) \rho_{\text{O}_2} + \dots$$

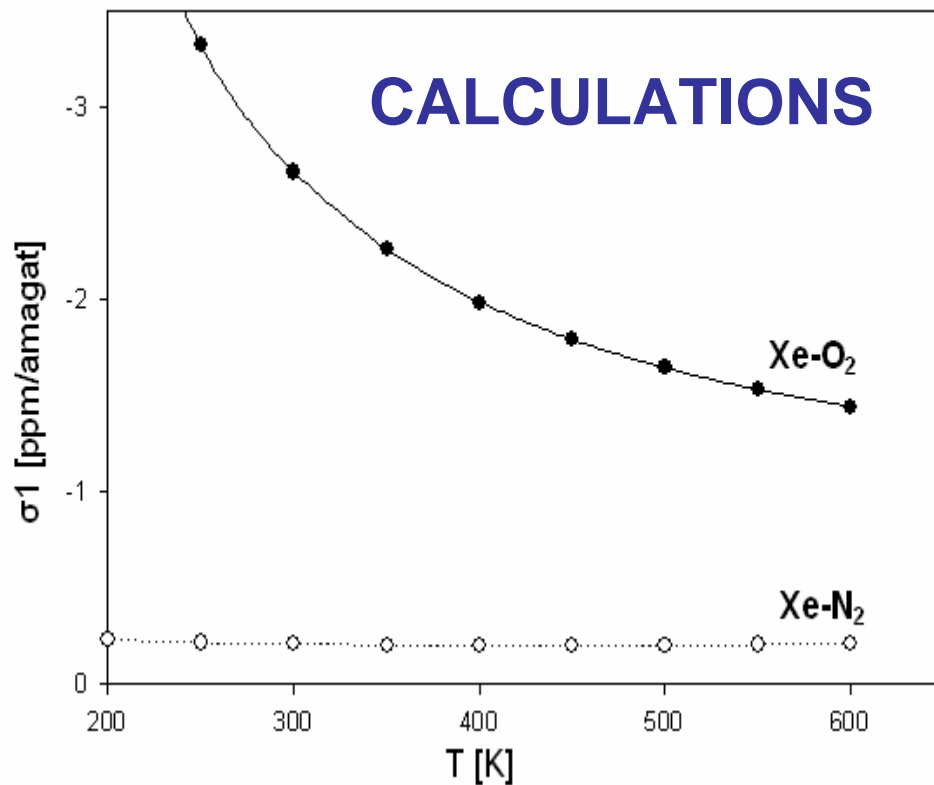
$$\sigma_1(\mathbf{T})_{\text{diam}} = 2\pi \int \{ \sigma(\mathbf{R}, \theta) - \sigma(\infty) \} e^{-V(\mathbf{R}, \theta)/kT} R^2 dR$$

$$\Delta \mathbf{B}/\mathbf{B}_0 = - \text{constant} * \int \{ \rho_{\text{spin}}(\mathbf{R}, \theta) - \rho_{\text{spin}}(\infty) \} e^{-V(\mathbf{R}, \theta)/kT} R^2 dR$$

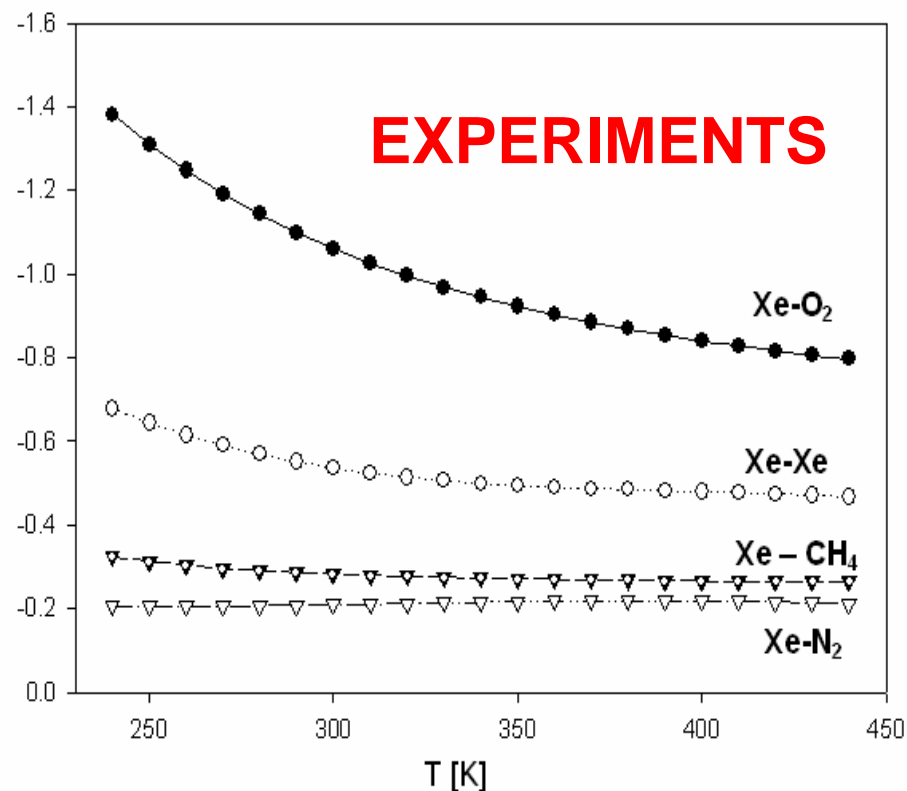
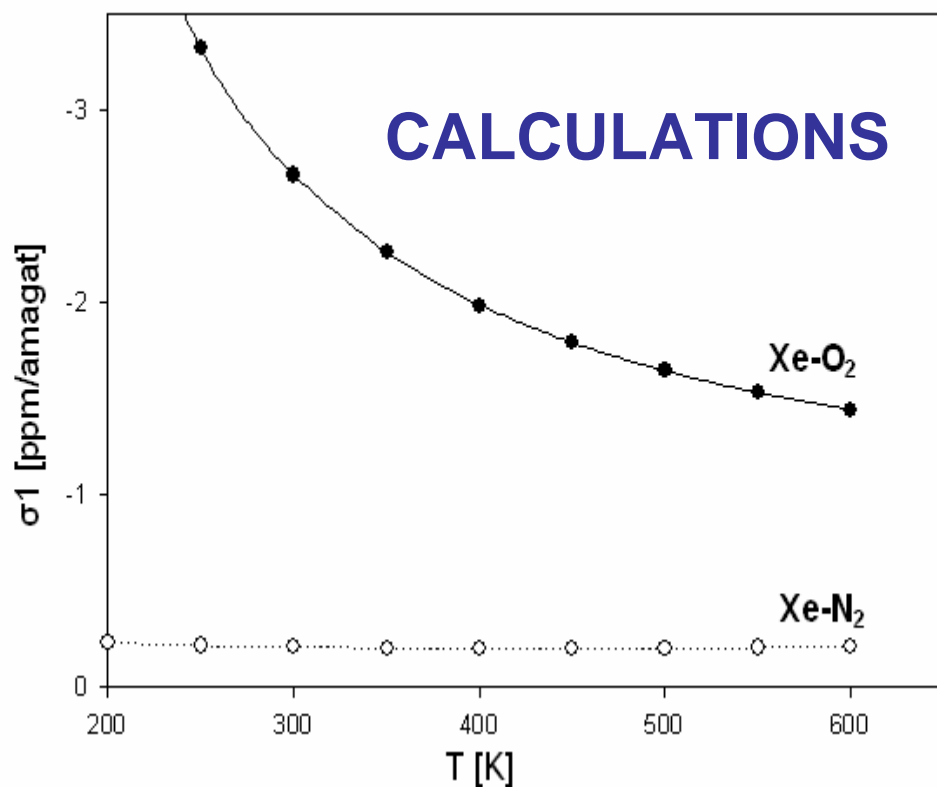


# Averaging over all Xe@O<sub>2</sub> configurations

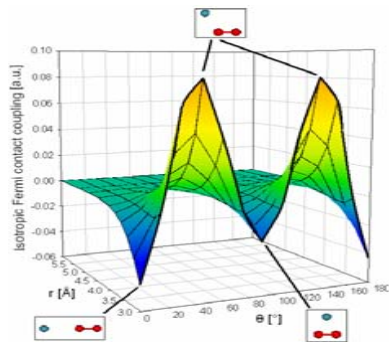
Averaging over all Xe@O<sub>2</sub> configurations, we obtain both (a) the usual chemical shift response to a diamagnetic molecule, and (b) also the Fermi contact contribution that only a paramagnetic molecule can provide.



# An experimental test:



Data from Jameson et al. Mol. Phys. 29, 1919 (1975), J. Chem. Phys. (1975-1978)



# CONCLUSIONS

**This analysis is a fundamental step towards a quantitative understanding of hyperfine shifts of Xe in the presence of paramagnetic centers in solids.**

- **magnitude of Xe chemical shift much larger than for diamagnetic counterpart**
- **obtained the correct sign**
- **correct temperature dependence**
- **our results suggest applications of Xe NMR as probe of paramagnetic centers in materials.**

THANKS to



and

Herbert Paaren Scholarship  
Lubrizol Scholarship

and to Devin N. Sears  
for help with Gaussian98