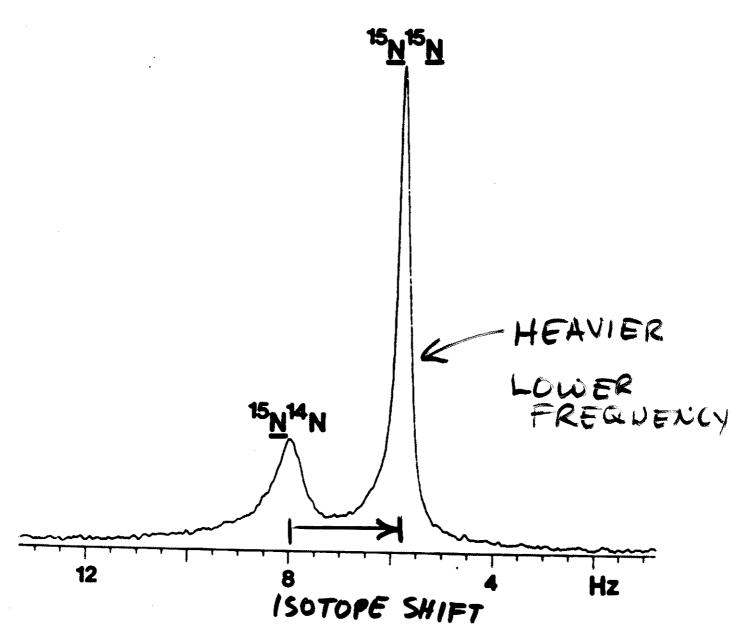
## Isotope Effects on NMR Chemical Shifts

Cynthia J. Jameson
University of Illinois at
Chicago

#### ISOTOPE EFFECTS ON NMR CHEMICAL SHIFTS

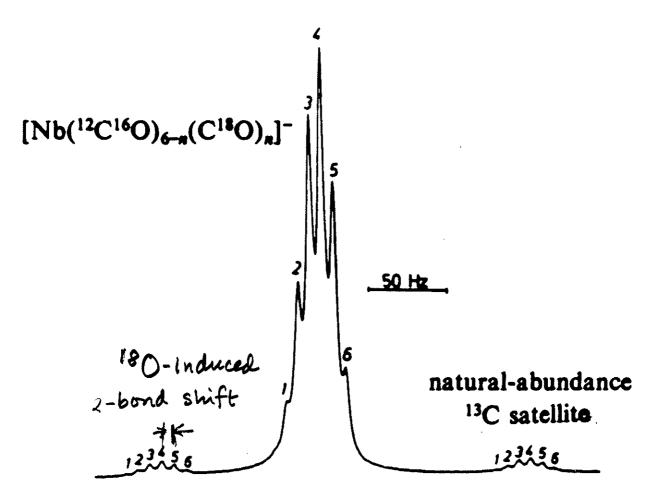
#### **OUTLINE**

- 1. Some empirical observations
- 2. Summary of empirical observations
- 3. Probability of finding a molecule at a given nuclear configuration
- 4. Dynamic averages: mass dependence
- 5. Nuclear shielding surfaces
  How does shielding change with increasing bond length and bond angle?
- 6. Trends in Electronic Factors
  - a. symmetry
  - b. net charge
  - c. number of lone pairs
  - d. bond order
- 7. Temperature dependence



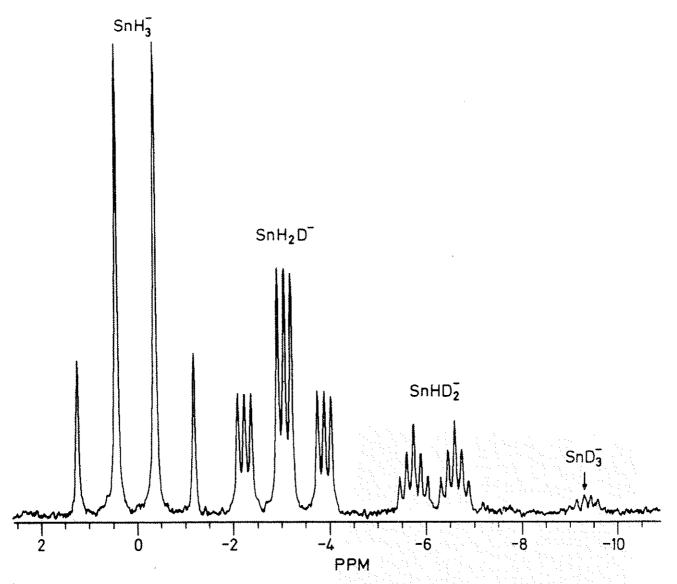
Nitrogen-15 NMR spectrum of <sup>14</sup>N<sup>15</sup>N and <sup>15</sup>N<sup>15</sup>N at 36.5977 MHz.

R.E. Wasylisher

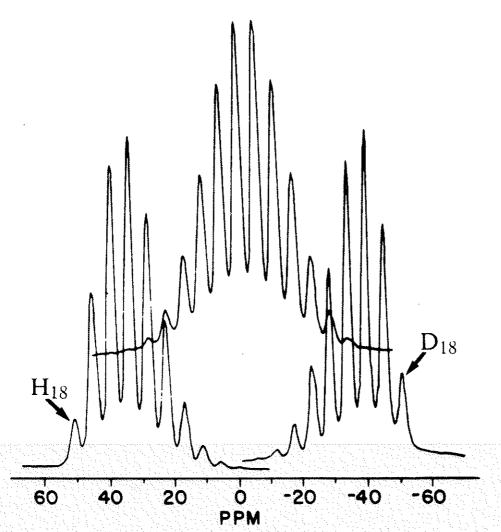


<sup>93</sup>Nb NMR spectra of [Nb(CO)<sub>6</sub>]<sup>-</sup> ion enriched in C<sup>18</sup>O taken at 88.29 MHz in acetone solution.

D. Rehdu



<sup>119</sup>Sn D-induced isotope shifts in the deuterated SnH<sub>3</sub> species are -3.281 ppm / D



 $^{59}$ Co NMR spectrum of 3 samples of deuterated  $^{59}$ Co[(NH<sub>3</sub>)<sub>6</sub>]<sup>2+</sup>

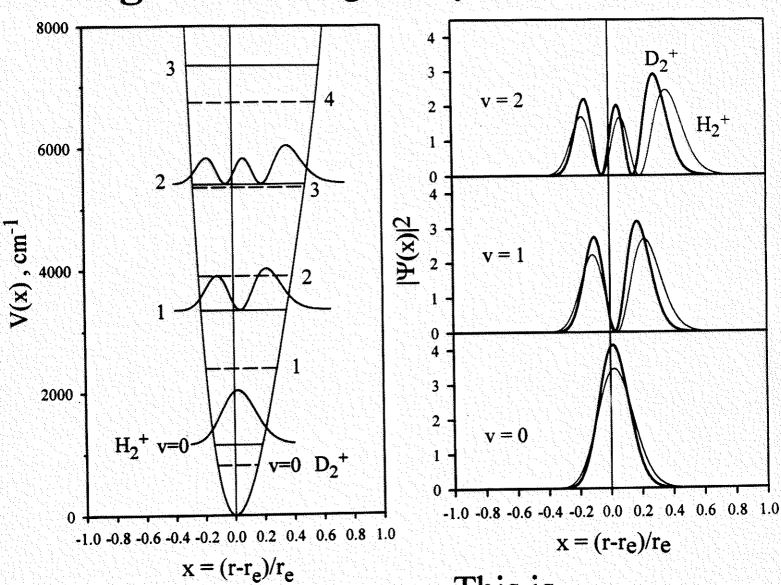
$$I = \frac{18!}{(1-d)^{p}(d)^{(18-p)}}$$
$$p!(18-p)!$$

d = overall deuterium fraction in sampled = 0.15, 0.50, 0.85 from left to rightRussell and Bryant

## Summary of empirical observations about isotope shifts in NMR:

- (1) Heavy isotopic substitution shifts the NMR signal of a nearby nucleus toward lower frequency (higher magnetic field).
- (2) The magnitude of the isotope shift is dependent on how remote the isotopic substitution is from the nucleus under observation.
- (3) The magnitude of the shift is a function of the resonant nucleus, reflecting the differences in the range of chemical shifts for the nuclei.
- (4) The magnitude of the shift is largest where the fractional change in mass upon isotopic substitution is largest. (largest for  $H \rightarrow D$  or T)
- (5) The magnitude of the shift is approximately proportional to the number of atoms in the molecule that have been substituted by isotopes.

# 3. The probability of finding a molecule at a given nuclear configuration is given by $|\Psi_{\text{vib}}|^2$ .



This is mass-dependent.

#### 4. Dynamic averages

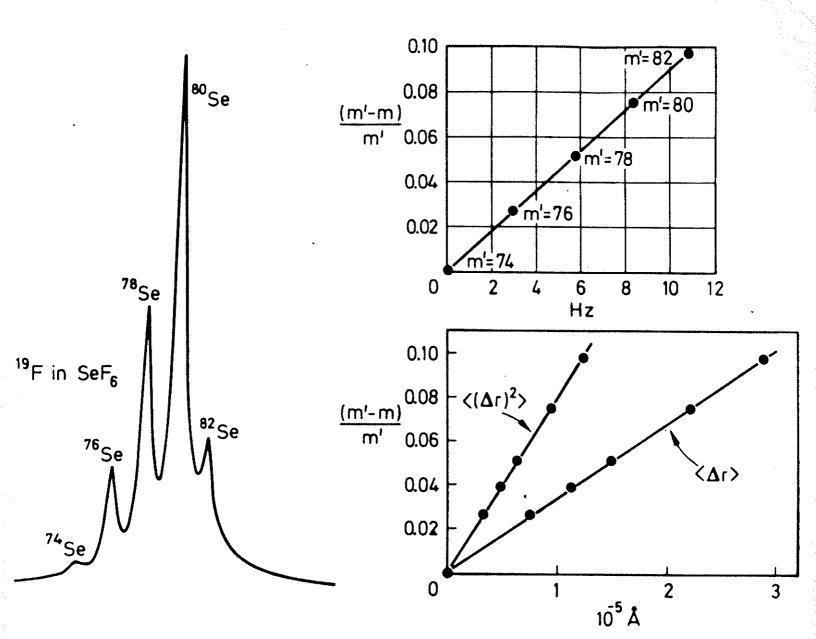
To get the average value of a molecular electronic property:

$$\langle P \rangle_{v} = \int_{-\infty}^{+\infty} |\Psi_{v}(x)|^{2} P(x) dx$$

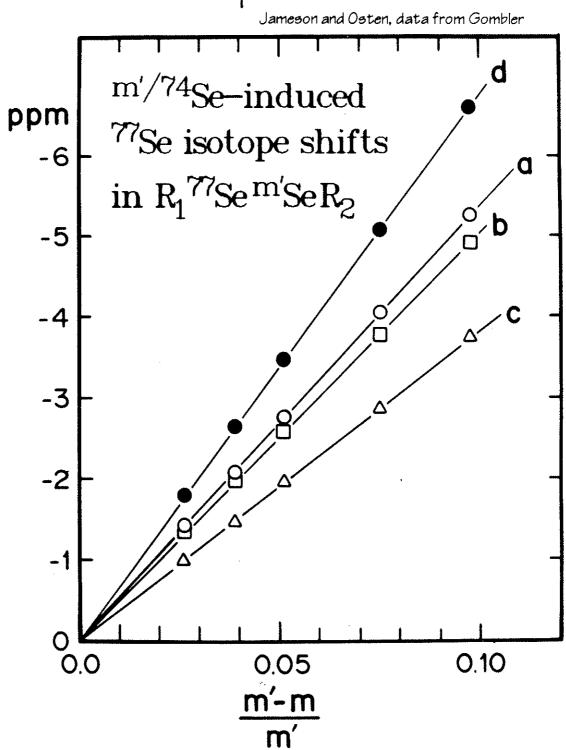
We can also find the rotational average. The thermal average is

$$\langle P \rangle^T = \frac{\sum\limits_{v,J,K} (2J+1)g_{Ns} \langle P \rangle_{vJK} \exp(-E_{vJK}/kT)}{\sum\limits_{v,J,K} (2J+1)g_{Ns} \exp(-E_{vJK}/kT)}$$

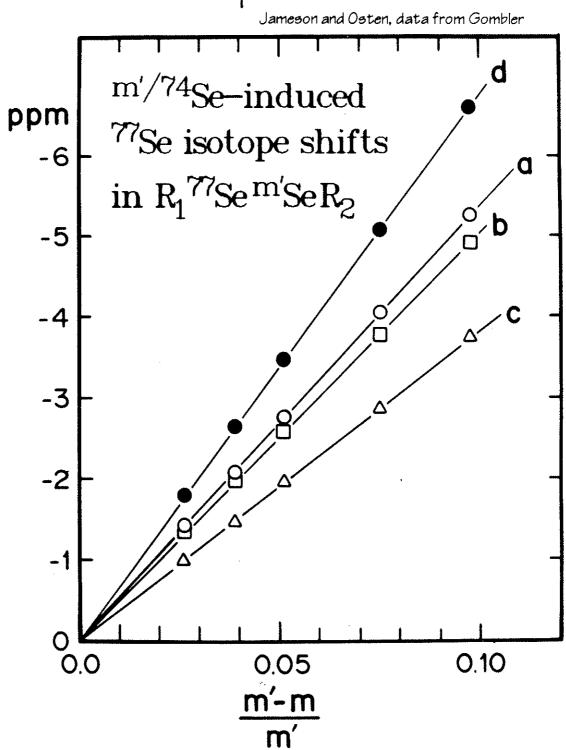
We can see that the low frequency vibrations are important, but also those vibrational modes that have large  $\langle P \rangle_{vJ}$ 

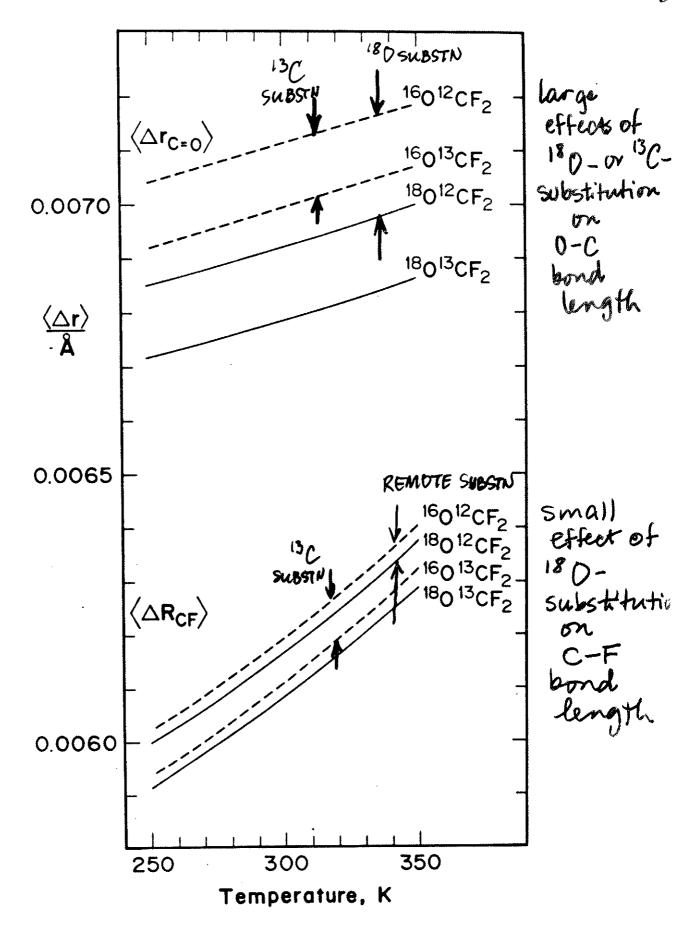


The mass-dependence of the dynamic part is predicted by theory: both  $\langle \Delta r \rangle$  and  $\langle (\Delta r)^2 \rangle$  depend on (m'-m)/m'. Here m and m' are the masses of the Se isotopes, 74 - 82.



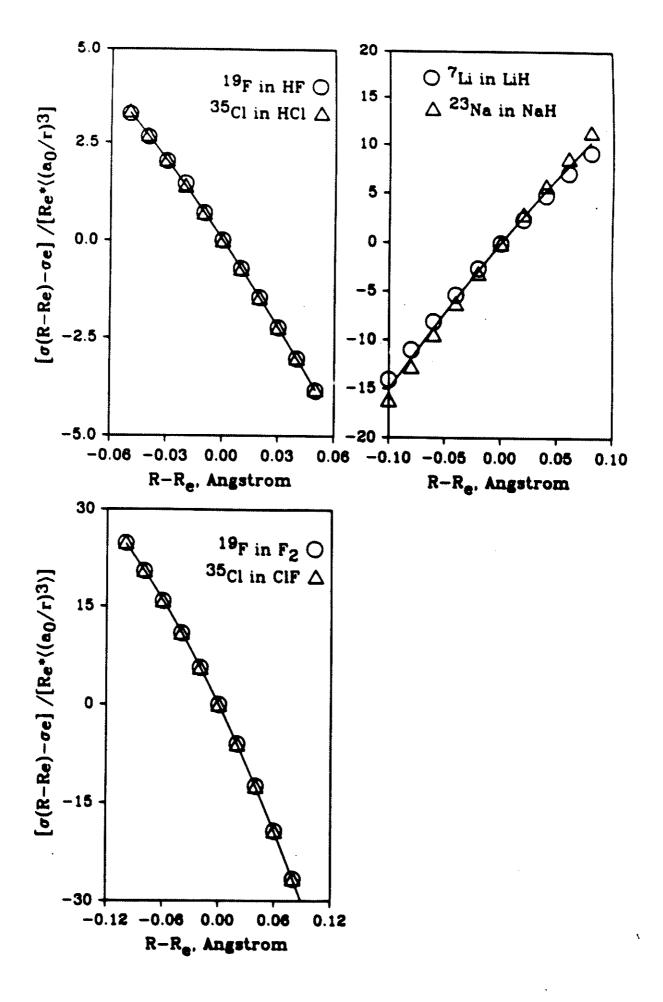
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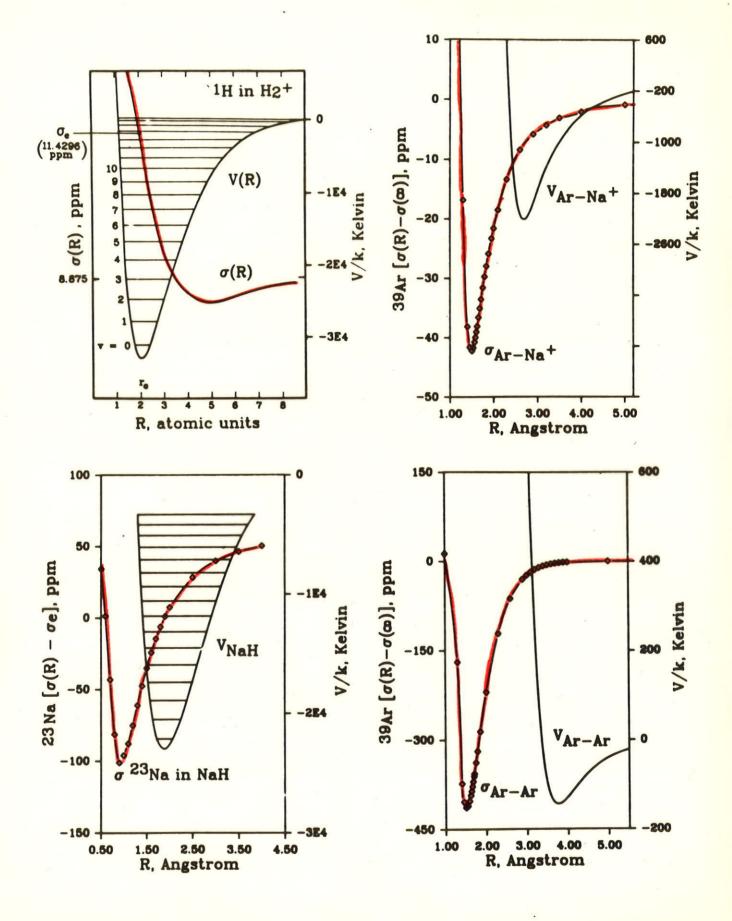


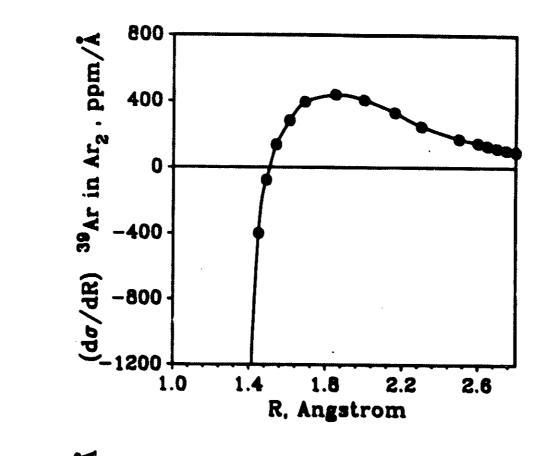


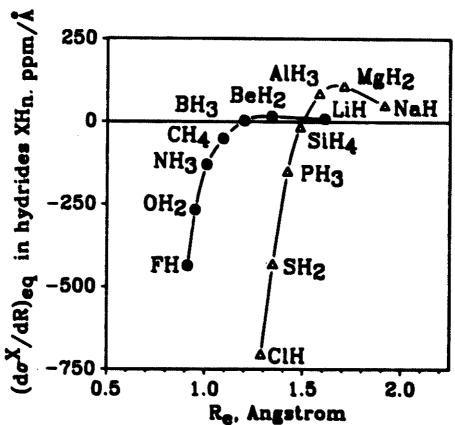
#### 5. NUCLEAR SHIELDING SURFACES

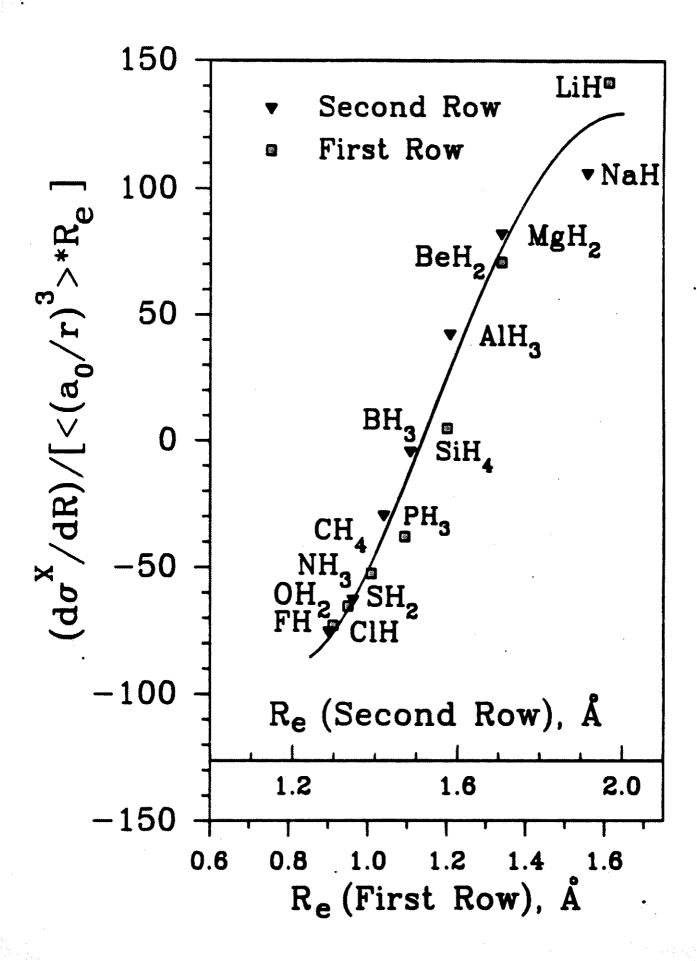
- a. How does shielding change with increasing bond length? variation of the derivative  $(\partial \sigma/\partial \mathbf{r})$
- b. What is  $(\partial \sigma/\partial \mathbf{r})_e$  i.e., the value at the equilibrium geometry of the molecule?
- c. How does  $(\partial \sigma/\partial \mathbf{r})_e$  change across the Periodic Table?
- d. When scaled, does the variation of  $(\partial \sigma/\partial \mathbf{r})_e$  across the Periodic Table make sense?
- e. Do the different signs of (∂σ/∂r)<sub>e</sub> make sense in terms of the variation of electron density with increasing internuclear distance?

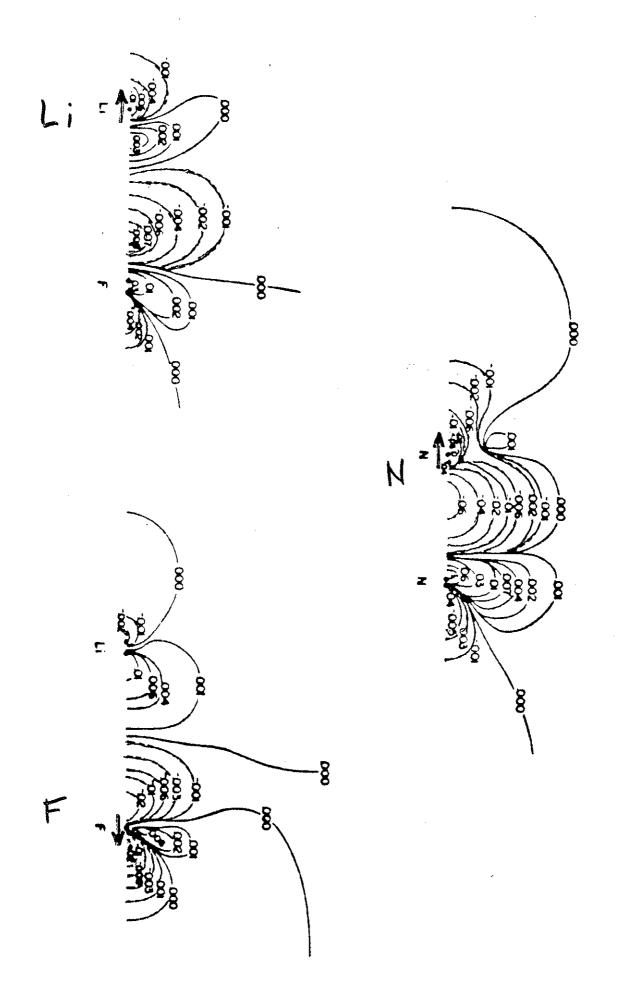








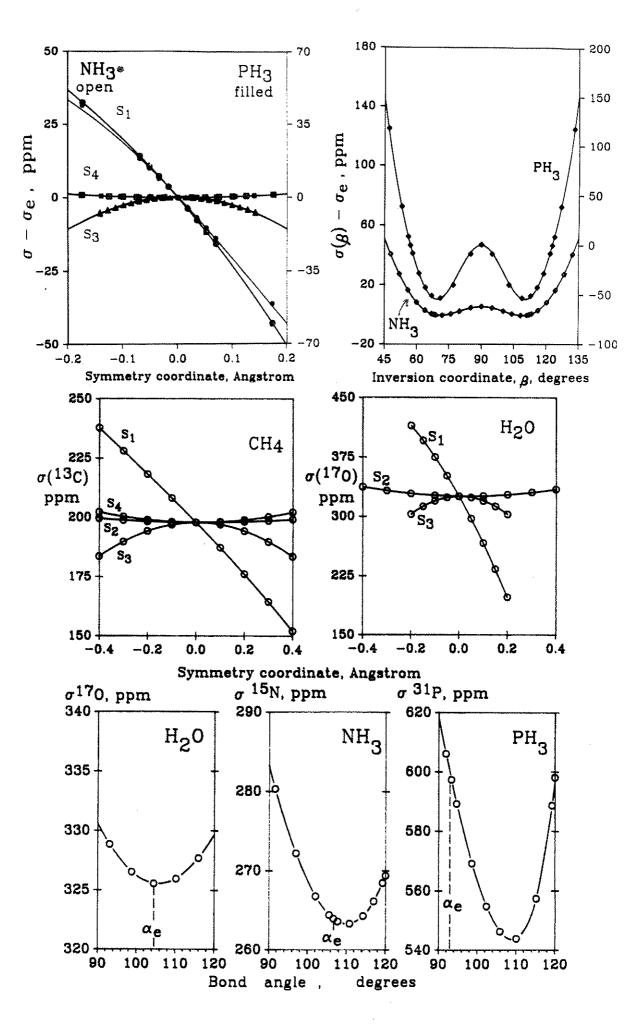




Electron density difference maps showing how electron density changes with bond extension

The charge redistribution in diatomic molecules, caused by extension of the bond:

- a) Charge is removed from along the internuclear axis and becomes concentrated in the perpendicular region. This leads to a quadrupolar character of the change in charge density, most pronounced for N and F. This leads to a larger paramagnetic contribution to the shielding, a net deshielding.
- b) On the other hand, charge is removed from the overlap region and builds up immediately in front of the *Li* nucleus in *LiF*. This leads to an increase in shielding.



#### 6. Trends in Electronic Factors

How does the shielding derivative  $(\partial \sigma^N/\partial r)$  (and thus also the isotope shift) change with

- a. position of N in the Periodic Table
- b. paramagnetic shielding of nucleus N in the chemical specie
- c. symmetry at site of nucleus N
- d. net charge of the chemical specie
- e. number of lone pairs on atom N
- f. bond order of the bonds to atom N
- g. change in pi density at aromatic site remote from atom N

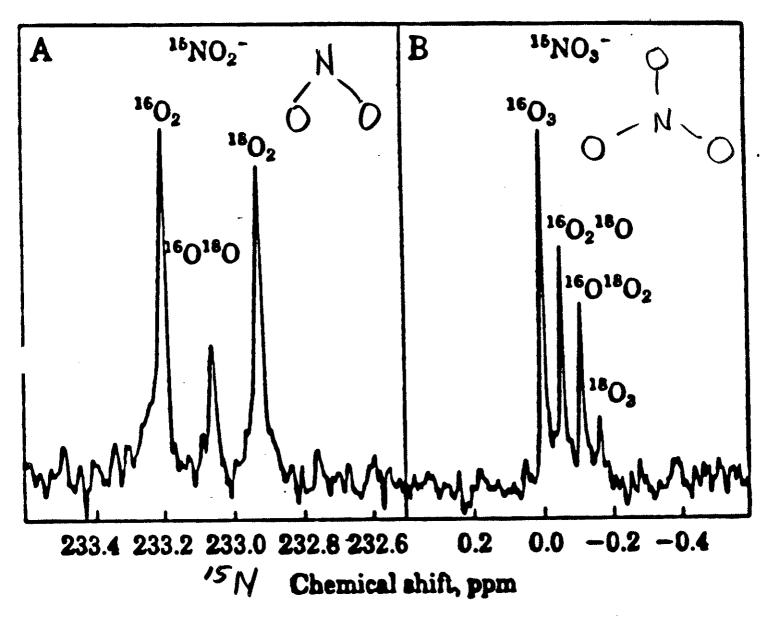


FIG. 1. <sup>18</sup>O/<sup>16</sup>O derivatives of nitrite and nitrate resolved by high-resolution <sup>16</sup>N NMR. (A) Solution of 100 mM [<sup>16</sup>N]nitrite in 0.25 M phosphate/0.5 mM carbonate/6% H<sub>2</sub><sup>18</sup>O, pH 7.5; 180 scans, 12-mm tube (4 ml). (B) Solution of [<sup>15</sup>N]nitrite in 0.25 M phosphate/5 mM rbonate/10% H<sub>2</sub><sup>18</sup>O, pH 7.0; 724 scans, 12-mm tube.

#### Effects of the net charge

<sup>119</sup> Sn in SnH <sub>4</sub> <sup>119</sup> Sn in SnH <sub>3</sub> +	$^{1}\Delta^{119}Sn(^{2/1}H)$ ppm per D -0.403 $-0.05 \pm 0.03$	emp. ( $\partial \sigma / \partial r$ ) ppm Å -1 -92 -10
	theor. $(\partial \sigma/\partial r)_e$	ppm Å <sup>-1</sup>
$^{11}B$ in $BH_3$	-3.5	
<sup>11</sup> B in BH <sub>4</sub> -	<b>-27</b> .	
<sup>27</sup> Al in AlH <sub>3</sub>	+84.2	
<sup>27</sup> Al in AlH <sub>4</sub> .	+11.6	3
	(∂ <b>o</b> /∂	Pr), ppm Å-1
	theor,	emp.
<sup>13</sup> C in HCN	-263	_
<sup>13</sup> C in CN	-538.7	-473
<sup>15</sup> N in HCN	-675.4	
<sup>15</sup> N in CN	-892.2	-872

		number of lone pairs	Average <sup>1</sup> Δ ppm per D
15N	in NH <sub>3</sub> NH <sub>4</sub> <sup>+</sup>	1 0	- 0.623 - 0.293
<sup>17</sup> O	in H <sub>2</sub> O H <sub>3</sub> O+	<b>2</b> 1	- 3.090 - 0.3
<sup>31</sup> P	in PH <sub>2</sub> · PH <sub>3</sub> PH <sub>4</sub> +	2 1 0	- 2.760 - 0.846 very small
<sup>119</sup> S	n in SnH <sub>3</sub> SnH <sub>4</sub> SnH <sub>3</sub> <sup>+</sup>	1 0 0	- 3.281 - 0.403 - 0.05 ± 0.03

#### Lone pair effects

	Ab ini	itio (30/3r).	ppm Å	From expt
N in	- 13	130.3 -144 (-70 - 67.9	) from LP)	-124 -60 , -65
P in	PH <sub>3</sub> -	150.8 - 154 (-4 - 52.9	7 from LP)	-585 -180 very small
Sn in	SnH <sub>3</sub> · SnH <sub>4</sub> SnH <sub>3</sub> ·	•		- 750 - 92 - 10

#### **EMPIRICAL TREND:**

greater bond order, larger isotope shifts

#### **THEORY:**

	(∂σ <sup>C</sup> /∂r <sub>CO</sub> ) <sub>e</sub> ppm Å <sup>-1</sup>		
	empirical	theoretical*	
CO	-456	-573.9	
O=C=O	-214, -220	-156.4	

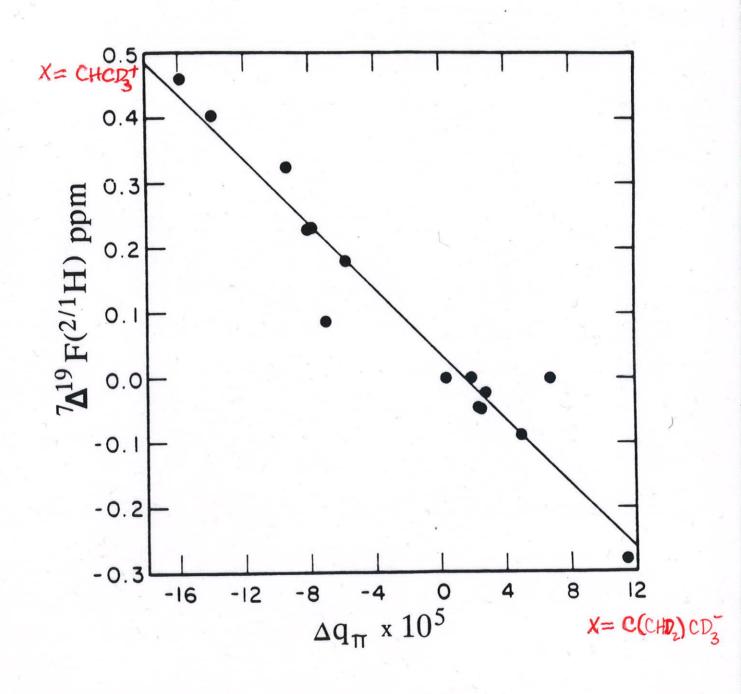
<sup>\*</sup> GIAO Chesnut

#### **SUMMARY OF TRENDS:**

	positive⁴	$(\partial \sigma/\partial r)_e \leftarrow$	- large negative
net charge	positive ion	neutral	negative ion
lone	0	1	2
pairs			
bond	1	2	3
order			

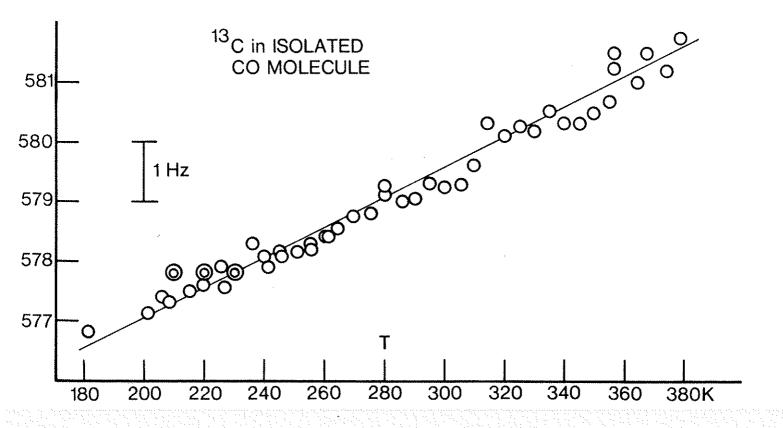
- DEPENDENCE of the MAGNITUDE of ISOTOPE SHIFT on the NMR MUCLEUS
- · THE SIGN OF the ISOTOPE SHIFT

19F D-induced isotope shifts in 4-fluorophenyl systems correlate with calculated change in pidensity induced by shortening the C-H(D) bond Forsyth and Yang

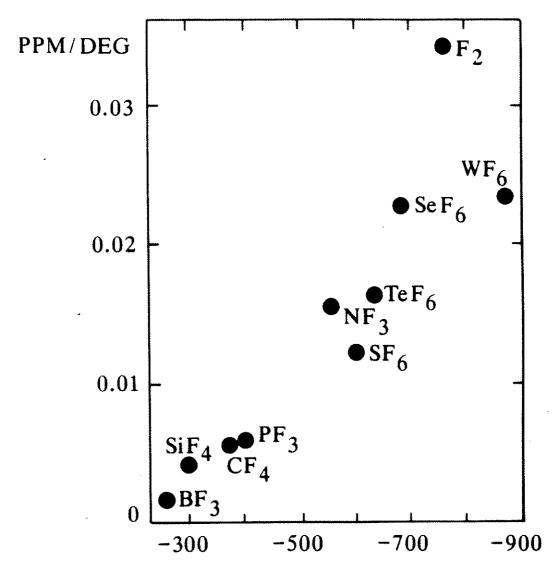


### 7. Temperature dependence of the chemical shift

- a. How the temperature dependence is related to isotope shifts
- b. How to calculate isotope shifts and their temperature dependence

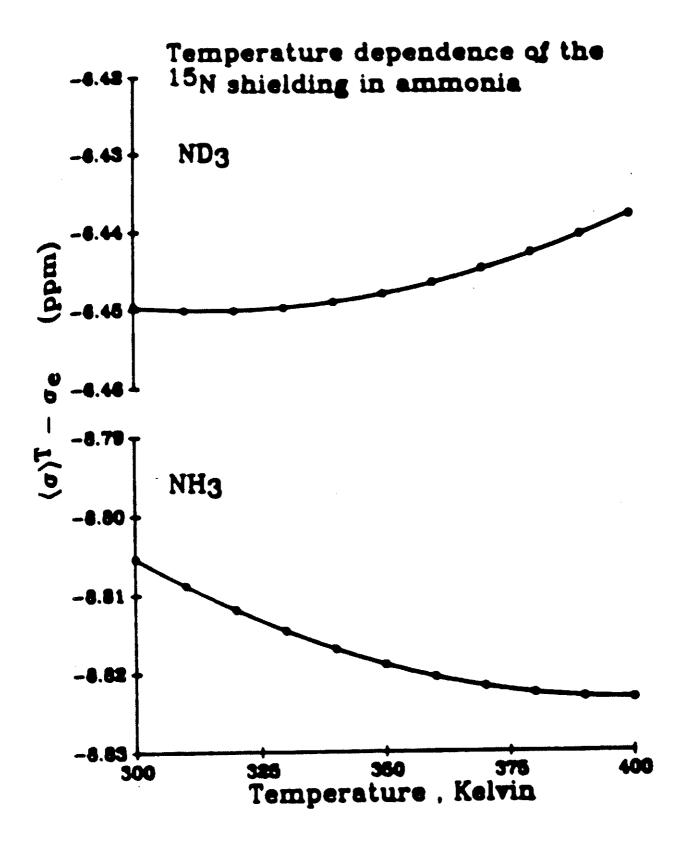


<sup>13</sup>C in isolated CO molecule becomes less shielded (moves to higher frequency) with increasing temperature. This is consistent with the sign and magnitude of the <sup>18</sup>O-induced isotope shift (lighter isotope, higher frequency).



PARAMAGNETIC SHIELDING σ<sup>P</sup>

The greater temperature dependence of the <sup>19</sup>F chemical shift goes with the more negative paramagnetic shielding. That is, most of the change in shielding upon bond extension comes from the change in the paramagnetic part.



POTENTIAL ENERGY SURFACE:

$$V = \frac{1}{2!} \sum_{i} \omega_{i} q_{i}^{2} + \frac{1}{3!} \sum_{ijk} \phi_{ijk} q_{i} q_{j} q_{k} + \cdots$$

PROPERTY SURFACE:

The particular, but could be any other molecular electronic property  $P = P_e + \sum_i P_i \, q_i + \sum_{ij} \frac{1}{2!} \, P_{ij} \, q_i \, q_j + \sum_{ijk} \frac{1}{3!} \, P_{ijk} \, q_i \, q_j \, q_k + \cdots$ 

To first order

$$\begin{split} \langle P \rangle_{n} &= \langle X_{n}^{(0)} + \lambda X_{n}^{(1)} | P | X_{n}^{(0)} + \lambda X_{n}^{(1)} \rangle \\ &= \langle X_{n}^{(0)} | P | X_{n}^{(0)} \rangle + 2\lambda \langle X_{n}^{(0)} | P | X_{n}^{(1)} \rangle \end{split}$$

VIBRATIONAL AVERAGE:  $\langle q_s q_s \rangle = (v_s + \frac{d_s}{2}) \delta_{ss}$ 

$$\langle q_{s'} \rangle = -\sum_{s} (\phi_{sss'}/2\omega_{s'})(v_s + \frac{d_s}{2})$$

non-vanishing only when s' is totally symmetric

$$+rac{1}{4\pi c\omega_{\mathrm{s}}} \left(rac{1}{\mathrm{hc\omega_{\mathrm{s}}}}
ight)^{1/2} \sum_{\alpha} rac{\mathrm{a}_{\mathrm{s'}}^{\alpha\alpha}}{\left(\mathrm{I}_{\alpha\alpha}^{(\mathrm{e})}
ight)^{2}} \langle J_{\alpha}^{2} \rangle$$
 rot. contrib.

$$\langle P \rangle = P_0 + \sum_{s} A_s (v_s + \frac{1}{2}) + \sum_{s \leq s'} B_{ss'} (v_s + \frac{1}{2}) (v_{s'} + \frac{1}{2})$$

where anharmonicity in the senty change of the constants 
$$A_{s} = P_{ss}/2 - \sum_{s'} P_{s'} \phi_{sss'}/2\omega_{s'}$$

$$P_0 = P_e + \sum_s \frac{1}{64} P_{ssss} - \sum_s 7 P_{sss} \phi_{sss} / 288 \omega_s$$

$$+ \sum_{s \neq s} 3 P_{sss} \phi_{sss} / \omega_s / 32 (4 \omega_s^2 - \omega_s^2)$$

$$B_{ss} = P_{ssss}/16 - \sum_{s'} P_{sss'} \phi_{sss'} (8\omega_s^2 - 3\omega_{s'}^2)/8\omega_{s'} (4\omega_s^2 - \omega_{s'}^2)$$

#### THE OBSERVED THERMAL AVERAGE:

$$\langle P \rangle^T = \frac{\sum\limits_{v,J,K} (2J+1)g_{Ns} \langle P \rangle_{vJK} exp(-E_{vJK}/kT)}{\sum\limits_{v,J,K} (2J+1)g_{Ns} exp(-E_{vJK}/kT)}$$

In terms of the curvilinear internal coordinates

$$P = P_{e} + \sum_{i} \left(\frac{\partial P}{\partial R_{i}}\right) R_{i} + \frac{1}{2!} \sum_{ij} \left(\frac{\partial^{2} P}{\partial R_{i} \partial R_{j}}\right)_{e} R_{i} R_{j}$$
$$+ \frac{1}{3!} \sum_{ijk} \left(\frac{\partial^{3} P}{\partial R_{i} \partial R_{j} \partial R_{k}}\right) R_{i} R_{j} R_{k} + \dots$$

For example, in CH<sub>4</sub>-type molecules

$$\langle \Delta \mathbf{r}_1 \rangle = \overline{\overline{L}}_1^{1} \langle \mathbf{q}_1 \rangle + \frac{1}{2} \overline{\overline{L}}_1^{2\mathbf{a}2\mathbf{a}} \langle \mathbf{q}_2^2 \rangle + \frac{1}{2} \overline{\overline{L}}_1^{3\mathbf{x}3\mathbf{x}} \langle \mathbf{q}_3^2 \rangle + \frac{1}{2} \overline{\overline{L}}_1^{4\mathbf{x}4\mathbf{x}} \langle \mathbf{q}_4^2 \rangle + \dots$$
 
$$\langle (\Delta \mathbf{r}_1)^2 \rangle = (\overline{\overline{L}}_1^1)^2 \langle \mathbf{q}_1^2 \rangle + (\overline{\overline{L}}_1^{3\mathbf{x}})^2 \langle \mathbf{q}_3^2 \rangle + (\overline{\overline{L}}_1^{4\mathbf{x}})^2 \langle \mathbf{q}_4^2 \rangle + \dots$$

#### **SUMMARY**

- Where an isotopic label is introduced in a molecule, every neighboring NMR nucleus experiences a slight chemical shift.
- If labeling is less than 100%, the resonant nuclei in both the labeled and the unlabeled molecules are observed, with intensities according to statistical distribution.
- The magnitude of the shift depends on the fractional mass change at the isotope substitution site, on the remoteness of the resonant nucleus from the substitution site, and on the sensitivity of the chemical shift of the resonant nucleus; the latter is reflected by the chemical shift range of that nucleus.

- The isotope shift is just one more powerful tool in which a very selective tag carries with it the same wealth of information as the chemical shift itself.
- Isotope shifts provide a more stringent test of ab initio calculations of chemical shifts in specific molecules, being directly related to the slopes on the mathematical surface that describes the variation of the chemical shift with the molecular geometry.
- In a more general sense, the trends in the thousands of isotope shifts that have been accumulated provide insight into the general nature of these chemical shift surfaces, in terms of the dependence of the details of the surface on the nature of the chemical bond, the net charge of the molecule, bond orders, presence of lone pairs, etc.

• With a method of estimating the dynamic factors in isotope shifts, we educe from tables of data on isotope shifts the general trends in the electronic factors (how the NMR chemical shifts depend on bond length for example), which have been recently verified by ab initio calculations.

#### **ACKNOWLEDGMENT**

Thanks to Prof. Poul Erik Hansen for organizing the June 1995 Summer School on Isotope Effects as a Tool in Basic and Environmental Science (UNESCO) at Roskilde University, Denmark





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