

Problem Set 3 Answers On Separation of Variables

1. (a) $\{-(\hbar^2/2M_1R^2)\partial^2/\partial\phi_1^2 - (\hbar^2/2M_2R^2)\partial^2/\partial\phi_2^2\} \Psi(\phi_1, \phi_2) = E\Psi(\phi_1, \phi_2)$

Let $\Psi(\phi_1, \phi_2) = \psi(\phi_1) \cdot \psi(\phi_2)$

$\{-(\hbar^2/2M_1R^2)\partial^2/\partial\phi_1^2 - (\hbar^2/2M_2R^2)\partial^2/\partial\phi_2^2\} \psi(\phi_1) \cdot \psi(\phi_2) = E\psi(\phi_1) \cdot \psi(\phi_2)$

$\psi(\phi_2)\{-(\hbar^2/2M_1R^2)\partial^2/\partial\phi_1^2 \psi(\phi_1)\} + \psi(\phi_1)\{-(\hbar^2/2M_2R^2)\partial^2/\partial\phi_2^2 \psi(\phi_2)\} = E\psi(\phi_1) \cdot \psi(\phi_2)$

Dividing both sides by $\psi(\phi_1) \cdot \psi(\phi_2)$

$$\frac{\{-(\hbar^2/2M_1R^2)\partial^2/\partial\phi_1^2 \psi(\phi_1)\}}{\psi(\phi_1)} + \frac{\{-(\hbar^2/2M_2R^2)\partial^2/\partial\phi_2^2 \psi(\phi_2)\}}{\psi(\phi_2)} = E$$

Since ϕ_1 and ϕ_2 are independent variables, then the first term has nothing to do with the second, so the only way the sum can equal a constant E is if the first term is equal to a constant E_1 and the second term is equal to a constant E_2 such that

$$E_1 + E_2 = E$$

Therefore

$$\frac{\{-(\hbar^2/2M_1R^2)\partial^2/\partial\phi_1^2 \psi(\phi_1)\}}{\psi(\phi_1)} = E_1$$

$$\frac{\{-(\hbar^2/2M_2R^2)\partial^2/\partial\phi_2^2 \psi(\phi_2)\}}{\psi(\phi_2)} = E_2$$

Solving,

$$\{-(\hbar^2/2M_1R^2)\partial^2/\partial\phi_1^2 \psi(\phi_1)\} = E_1\psi(\phi_1)$$

we find $\psi_{k_1}(\phi_1) = (1/\sqrt{2\pi}) \exp[i k_1 \phi_1]$ satisfies this equation with eigenvalues $k_1^2 \hbar^2/2M_1R^2$

Therefore

$$\Psi(\phi_1, \phi_2) = (1/\sqrt{2\pi}) \exp[i k_1 \phi_1] \cdot (1/\sqrt{2\pi}) \exp[i k_2 \phi_2].$$

$$E = k_1^2 \hbar^2/2M_1R^2 + k_2^2 \hbar^2/2M_2R^2$$

2. Similarly for 6 electrons:

(a) $\{-(\hbar^2/2m_eR^2)\partial^2/\partial\phi_1^2 - (\hbar^2/2M_2R^2)\partial^2/\partial\phi_2^2 - (\hbar^2/2m_eR^2)\partial^2/\partial\phi_3^2 - (\hbar^2/2m_eR^2)\partial^2/\partial\phi_4^2 - (\hbar^2/2m_eR^2)\partial^2/\partial\phi_5^2 - (\hbar^2/2m_eR^2)\partial^2/\partial\phi_6^2\} \Psi(\phi_1, \phi_2, \phi_3, \phi_4, \phi_5, \phi_6) = E\Psi(\phi_1, \phi_2, \phi_3, \phi_4, \phi_5, \phi_6)$

(b) $\Psi(\phi_1, \phi_2, \phi_3, \phi_4, \phi_5, \phi_6) = \psi(\phi_1) \cdot \psi(\phi_2) \cdot \psi(\phi_3) \cdot \psi(\phi_4) \cdot \psi(\phi_5) \cdot \psi(\phi_6)$

$$\Psi(\phi_1, \phi_2, \phi_3, \phi_4, \phi_5, \phi_6) = (1/\sqrt{2\pi}) \exp[i k_1 \phi_1] \cdot (1/\sqrt{2\pi}) \exp[i k_2 \phi_2] \cdot (1/\sqrt{2\pi}) \exp[i k_3 \phi_3] \cdot (1/\sqrt{2\pi}) \exp[i k_4 \phi_4] \cdot (1/\sqrt{2\pi}) \exp[i k_5 \phi_5] \cdot (1/\sqrt{2\pi}) \exp[i k_6 \phi_6]$$

$$E = k_1^2 \hbar^2/2m_eR^2 + k_2^2 \hbar^2/2m_eR^2 + k_3^2 \hbar^2/2m_eR^2 + k_4^2 \hbar^2/2m_eR^2 + k_5^2 \hbar^2/2m_eR^2 + k_6^2 \hbar^2/2m_eR^2$$

where $k_1 = 0, \pm 1, \pm 2, \dots$, similarly for k_2 , etc.

(c) Invoking the condition that no two electrons can have the same k value, the lowest energy can be $(k_1 k_2 k_3 k_4 k_5 k_6) = (0011-1-1)$ with total energy given by $(0+0+1+1+1+1) \hbar^2/2m_e R^2 = 4 \hbar^2/2m_e R^2$

The next lowest energy will be given by $(k_1 k_2 k_3 k_4 k_5 k_6) = (0011-12)$ or $(0011-1-2)$ or $(0012-1-1)$ or $(001-2-1-1)$ [not including permuting the electron's numbers] with energy $(0+0+1+1+1+4) \hbar^2/2m_e R^2 = 7 \hbar^2/2m_e R^2$

The next lowest energy will be given by $(k_1 k_2 k_3 k_4 k_5 k_6) = (011-1-12)$ or $(011-1-1-2)$ with energy $(0+1+1+1+1+4) \hbar^2/2m_e R^2 = 8 \hbar^2/2m_e R^2$

(d) $\int \psi_{k'}^* \times \psi_{k'} d\tau$ where $x = R \cos \phi =$

The peak intensity depends on the square of the integral:

$$\iiint \Psi^*(\phi_1, \phi_2, \phi_3, \phi_4, \phi_5, \phi_6) R [\cos \phi_1 + \cos \phi_2 + \cos \phi_3 + \cos \phi_4 + \cos \phi_5 + \cos \phi_6] \Psi'(\phi_1, \phi_2, \phi_3, \phi_4, \phi_5, \phi_6) d\phi_1 d\phi_2 d\phi_3 d\phi_4 d\phi_5 d\phi_6$$

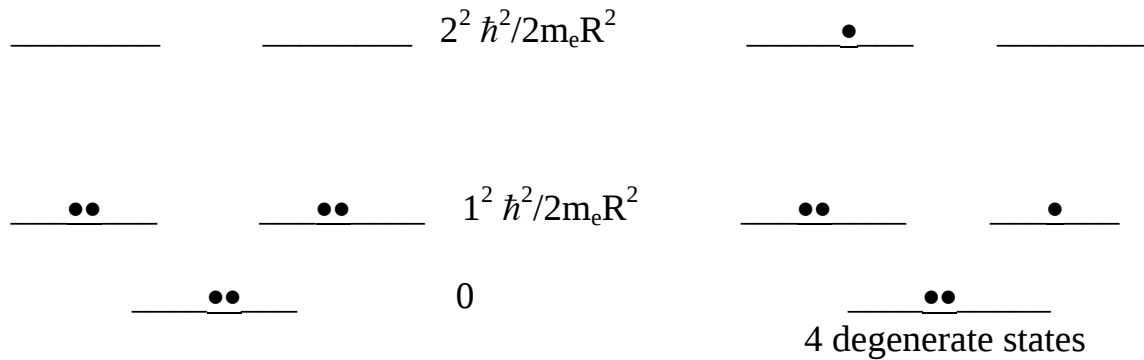
This integral leads to a sum of integrals:

$$\begin{aligned} & \delta(k_2, k_2') \bullet \delta(k_3, k_3') \bullet \delta(k_4, k_4') \bullet \delta(k_5, k_5') \bullet \delta(k_6, k_6') \bullet \int \psi_{k_1}^*(\phi_1) R \cos \phi_1 \psi'_{k_1'}(\phi_1) d\phi_1 \\ & + \delta(k_1, k_1') \bullet \delta(k_3, k_3') \bullet \delta(k_4, k_4') \bullet \delta(k_5, k_5') \bullet \delta(k_6, k_6') \bullet \int \psi_{k_2}^*(\phi_2) R \cos \phi_2 \psi'_{k_2'}(\phi_2) d\phi_2 \\ & + \delta(k_1, k_1') \bullet \delta(k_2, k_2') \bullet \delta(k_4, k_4') \bullet \delta(k_5, k_5') \bullet \delta(k_6, k_6') \bullet \int \psi_{k_3}^*(\phi_3) R \cos \phi_3 \psi'_{k_3'}(\phi_3) d\phi_3 \\ & + \delta(k_1, k_1') \bullet \delta(k_2, k_2') \bullet \delta(k_3, k_3') \bullet \delta(k_5, k_5') \bullet \delta(k_6, k_6') \bullet \int \psi_{k_4}^*(\phi_4) R \cos \phi_4 \psi'_{k_4'}(\phi_4) d\phi_4 \\ & + \delta(k_1, k_1') \bullet \delta(k_2, k_2') \bullet \delta(k_3, k_3') \bullet \delta(k_4, k_4') \bullet \delta(k_6, k_6') \bullet \int \psi_{k_5}^*(\phi_5) R \cos \phi_5 \psi'_{k_5'}(\phi_5) d\phi_5 \\ & + \delta(k_1, k_1') \bullet \delta(k_2, k_2') \bullet \delta(k_3, k_3') \bullet \delta(k_4, k_4') \bullet \delta(k_5, k_5') \bullet \int \psi_{k_6}^*(\phi_6) R \cos \phi_6 \psi'_{k_6'}(\phi_6) d\phi_6 \end{aligned}$$

where the delta function $\delta(k_2, k_2') = 1$ if $k_2 = k_2'$ and 0 otherwise. These delta functions result from the fact that the eigenfunctions are orthonormal.

The only term that can survive is when $k_i = k_i'$ for all except that corresponding to the operator $R \cos \phi_j$. This means that a transition is only possible between 6-electron states $\Psi(\phi_1, \phi_2, \phi_3, \phi_4, \phi_5, \phi_6)$ that differ by no more than one electron's quantum number, and it is that one for which the integral $\int \psi_{k_j}^*(\phi_j) R \cos \phi_j \psi'_{k_j'}(\phi_j) d\phi_j$ has to be evaluated. In turn, this latter integral is evaluated by using the identity $\cos \phi = \frac{1}{2} [\exp(i\phi) + \exp(-i\phi)]$

$$\begin{aligned} & \int (1/\sqrt{2\pi}) \exp[-i k_1 \phi_1] \bullet \frac{1}{2} [\exp(i\phi_1) + \exp(-i\phi_1)] \bullet (1/\sqrt{2\pi}) \exp[i k_1' \phi_1] d\phi_1 = \\ & (1/4\pi) [\delta(k_1, k_1' + 1) + \delta(k_1, k_1' - 1)], \quad \text{that is } k_1 \text{ and } k_1' \text{ have to differ by 1.} \end{aligned}$$



Total energy = $(4) \hbar^2/2m_e R^2$

Total energy = $(7) \hbar^2/2m_e R^2$

(e) perimeter of the circle is $6(1.394) \text{ \AA} = 2\pi R$; $R = 1.331 \text{ \AA}$. The longest wavelength corresponds to transition from ground $(4)\hbar^2/2m_e R^2$ to excited state $(7)\hbar^2/2m_e R^2$ with transition energy = $3 \hbar^2/2m_e R^2$. Convert this energy to wavelength by $\Delta E = hc/\lambda$ or $\lambda = (8\pi^2 m_e R^2 c/3h)$

electron mass = $9.10938188 \times 10^{-28} \text{ g}$, $R = 1.331 \text{ \AA}$, $c = 3 \times 10^{10} \text{ cm s}^{-1}$

$$\lambda = \frac{8\pi^2 9.10938188 \times 10^{-28} (1.331 \times 10^{-8})^2 3 \times 10^{10}}{3 \times 6.62618 \times 10^{-27}}$$

$\lambda = 1920 \times 10^{-8} \text{ cm}$

(f) Two condensed rings outer perimeter = $10(1.394) \text{ \AA} = 2\pi R$.

For N_r rings, $[6+4(N_r - 1)](1.394) = 2\pi R$;

The radius of the equivalent "circle" is $R = [6+4(N_r - 1)](1.394)/2\pi$

The number of electrons is $6+4(N_r - 1)$.

The ground state energy is $4+4(2^2) + 4(3^2) + 4(4^2) + \dots$

For $N_r=2$: $E_{\text{grd}} = [4+4(2^2)] \hbar^2/2m_e (10 \times 1.394/2\pi)^2$

For $N_r=3$: $E_{\text{grd}} = [4+4(2^2) + 4(3^2)] \hbar^2/2m_e (14 \times 1.394/2\pi)^2$

For $N_r=4$: $E_{\text{grd}} = [4+4(2^2) + 4(3^2) + 4(4^2)] \hbar^2/2m_e (18 \times 1.394/2\pi)^2$

For the lowest excited states:

For $N_r=2$: $E_{\text{excited}} = [4+3(2^2) + 3^2] \hbar^2/2m_e (10 \times 1.394/2\pi)^2$

For $N_r=3$: $E_{\text{excited}} = [4+4(2^2) + 3(3^2) + 4^2] \hbar^2/2m_e (14 \times 1.394/2\pi)^2$

For $N_r=4$: $E_{\text{excited}} = [4+4(2^2) + 4(3^2) + 3(4^2) + 5^2] \hbar^2/2m_e (18 \times 1.394/2\pi)^2$

The energy of the transition is

For $N_r=2$: $\Delta E = [3^2 - 2^2] \hbar^2/2m_e (10 \times 1.394/2\pi)^2$

For $N_r=3$: $\Delta E = [4^2 - 3^2] \hbar^2/2m_e (14 \times 1.394/2\pi)^2$

For $N_r=4$: $\Delta E = [5^2 - 4^2] \hbar^2/2m_e (18 \times 1.394/2\pi)^2$

$$\Delta E = hc/\lambda$$

$$\lambda = (2m_e 1.394^2 c/h)(10)^2 / [3^2 - 2^2] \text{ for } N_r = 2$$

$$\lambda = (2m_e 1.394^2 c/h)(14)^2 / [4^2 - 3^2] \text{ for } N_r = 3$$

$$\lambda = (2m_e 1.394^2 c/h)(18)^2 / [5^2 - 4^2] \text{ for } N_r = 4$$

$$(2m_e 1.394^2 c/h) = \frac{2 \times 9.10938188 \times 10^{-28} (1.394 \times 10^{-8})^2 \times 3 \times 10^{10}}{6.62618 \times 10^{-27}} = 16.03 \times 10^{-7}$$

$$\lambda = 16.03 \times 10^{-7} (6)^2 / [2^2 - 1^2] = 16.03 \times 10^{-7} \times 12 = 1920 \text{ \AA for } N_r = 1$$

$$\lambda = 16.03 \times 10^{-7} (10)^2 / [3^2 - 2^2] = 16.03 \times 10^{-7} \times 20 = 3210 \text{ \AA for } N_r = 2$$

$$\lambda = 16.03 \times 10^{-7} (14)^2 / [4^2 - 3^2] = 16.03 \times 10^{-7} \times 28 = 4490 \text{ \AA for } N_r = 3$$

$$\lambda = 16.03 \times 10^{-7} (18)^2 / [5^2 - 4^2] = 16.03 \times 10^{-7} \times 36 = 5770 \text{ \AA for } N_r = 4$$

The number of rings determine the radius of the “circle” and also the number of electrons. The net result is $\lambda = 160.3 [12+8(N_r-1)] \text{ \AA}$

3.

$$x_{mn} = \frac{4a mn \{ (-1)^{\frac{m-n}{2}} - 1 \}}{\pi^2 \{ m^2 - n^2 \}^2} \text{ for } m \neq n$$

m\ n	1	2	3	4	5	6	7	...
1	////	2/9	0	4/225	0	6/1225	0	...
2	2/9	////	6/25	0	10/441	0	14/2025	...
3	0	6/25	////	12/49	0	18/729	0	...
4	4/225	0	12/49	////	20/81	0	28/1089	...
5 $-(8a/\pi^2)$	0	10/441	0	20/81	////	30/121	0	...
6	6/1225	0	18/729	0	30/121	////	42/169	...
7	0	14/2025	0	28/1089	0	42/169	////	...
8	8/3969	0	24/3025	0	40/1521	0	56/225	...
⋮	⋮	⋮	⋮	⋮	⋮	⋮	⋮	

//// although the integral is non-zero, when initial and final states are the same it does not correspond to a transition, instead an expectation value of x, $\langle x \rangle = a/2$.

$$\textcircled{3}(a) \quad x_{mn} = \int_0^a \Psi_m^*(x) \times \Psi(x) dx$$

$$= \frac{2}{a} \int_0^a \sin\left(\frac{m\pi}{a}x\right) \times \sin\left(\frac{n\pi}{a}x\right) dx$$

$$= \frac{1}{a} \int_0^a \left\{ \begin{array}{l} x \cos\left[(m-n)\frac{\pi}{a}x\right] \\ -x \cos\left[(m+n)\frac{\pi}{a}x\right] \end{array} \right\} dx$$

$$= \frac{1}{a} \left[\begin{array}{l} x \sin(m-n)\frac{\pi}{a}x + \frac{\cos(m-n)\frac{\pi}{a}x}{(m-n)^2\pi^2/a^2} \\ -x \sin(m+n)\frac{\pi}{a}x - \frac{\cos(m+n)\frac{\pi}{a}x}{(m+n)^2\pi^2/a^2} \end{array} \right]_0^a$$

$\sin$ is zero at both ends so we have:

$$= \frac{1}{a} \left\{ \frac{\cos(m-n)\pi - 1}{(m-n)^2\pi^2/a^2} - \frac{\cos(m+n)\pi - 1}{(m+n)^2\pi^2/a^2} \right\}$$

For $m=2$ $n=1$

$$x_{21} = \frac{1}{a} \left(\frac{a^2(\cos 2\pi - 1)}{1^2\pi^2} - \frac{a^2(\cos 3\pi - 1)}{3^2\pi^2} \right)$$

$$= \frac{a}{\pi^2} \left(\frac{-2}{1^2} - \frac{-2}{3^2} \right) = \frac{-8a}{\pi^2} \left(\frac{2}{9} \right)$$

$$(b) H = -\frac{\hbar^2}{2m} \left(\frac{d^2}{dx_1^2} + \frac{d^2}{dx_2^2} + \frac{d^2}{dx_3^2} + \frac{d^2}{dx_4^2} \right)$$

for butadiene $\rightarrow 1=2=3=4$

$$a = 5(1.40 \text{ \AA}) = 7.00 \text{ \AA} \quad (\text{given})$$

$$H(x_1, x_2, x_3, x_4) \Psi(x_1, x_2, x_3, x_4) = E \Psi(x_1, x_2, x_3, x_4)$$

Use separation of variables:

$$\text{Let } \Psi(x_1, x_2, x_3, x_4) = \psi(x_1) \cdot \psi(x_2) \cdot \psi(x_3) \cdot \psi(x_4)$$

$$\left\{ H(x_1) + H(x_2) + H(x_3) + H(x_4) \right\} \psi(x_1) \cdot \psi(x_2) \cdot \psi(x_3) \cdot \psi(x_4) = E \psi(x_1) \cdot \psi(x_2) \cdot \psi(x_3) \cdot \psi(x_4)$$

$$\frac{H(x_1)\psi(x_1)}{\psi(x_1)} + \frac{H(x_2)\psi(x_2)}{\psi(x_2)} + \frac{H(x_3)\psi(x_3)}{\psi(x_3)} + \frac{H(x_4)\psi(x_4)}{\psi(x_4)} = E$$

Therefore, each of the terms must be equal to some constant such that the sum of the four different constants sum to E . But we already know the answers to

$$H(x_1)\psi(x_1) = E_{n_1} \psi(x_1)$$

$$-\frac{\hbar^2}{2m} \frac{d^2}{dx_1^2} \psi(x_1) = E_{n_1} \psi(x_1) \text{ has solutions}$$

$$\psi(x_1) = \sqrt{\frac{2}{a}} \sin\left(\frac{n_1 \pi}{a} x_1\right)$$

$$E_{n_1} = \frac{n_1^2 \hbar^2}{8ma^2}$$

The others are of the same form,

$$\Psi(x_1, x_2, x_3, x_4) = \sqrt{\frac{2}{a}} \sin \frac{n_1 \pi}{a} x_1 \cdot \sqrt{\frac{2}{a}} \sin \frac{n_2 \pi}{a} x_2 \cdot \sqrt{\frac{2}{a}} \sin \frac{n_3 \pi}{a} x_3 \cdot \sqrt{\frac{2}{a}} \sin \frac{n_4 \pi}{a} x_4$$

$$\text{and } E = \frac{n_1^2 \hbar^2}{8ma^2} + \frac{n_2^2 \hbar^2}{8ma^2} + \frac{n_3^2 \hbar^2}{8ma^2} + \frac{n_4^2 \hbar^2}{8ma^2}$$

For particle 1 we have: $\frac{4^2 h^2}{8ma^2}$ Similarly for particles 2, 3, and 4

$$\frac{3^2 h^2}{8ma^2}$$

$$\frac{2^2 h^2}{8ma^2}$$

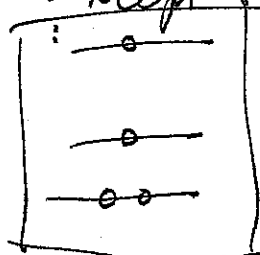
$$\frac{1^2 h^2}{8ma^2}$$

These are the one-electron states

Next energy levels:

$$n_1=1 \quad n_2=1 \quad n_3=2 \quad n_4=3 \text{ etc.}$$

(can assign all possible n values to each except that no more than ^{than} two can have the same n value.)



and so on.

$$E_{\text{excited}} = [2(1^2) + 1(2^2) + 1(3^2)] h^2 / 8ma^2 \text{ and so on.}$$

(c) The transitions depend on squaring the integral:

$$\int_0^a \dots \int_0^a \underbrace{\psi_{n_1}(x_1) \psi_{n_2}(x_2) \psi_{n_3}(x_3) \psi_{n_4}(x_4)}_{\text{initial state}} \underbrace{\psi_{n'_1}(x_1) \psi_{n'_2}(x_2) \psi_{n'_3}(x_3) \psi_{n'_4}(x_4)}_{\text{final state}} dx_1 \dots dx_4$$

Ground: 1122

1st excited 1123 etc.

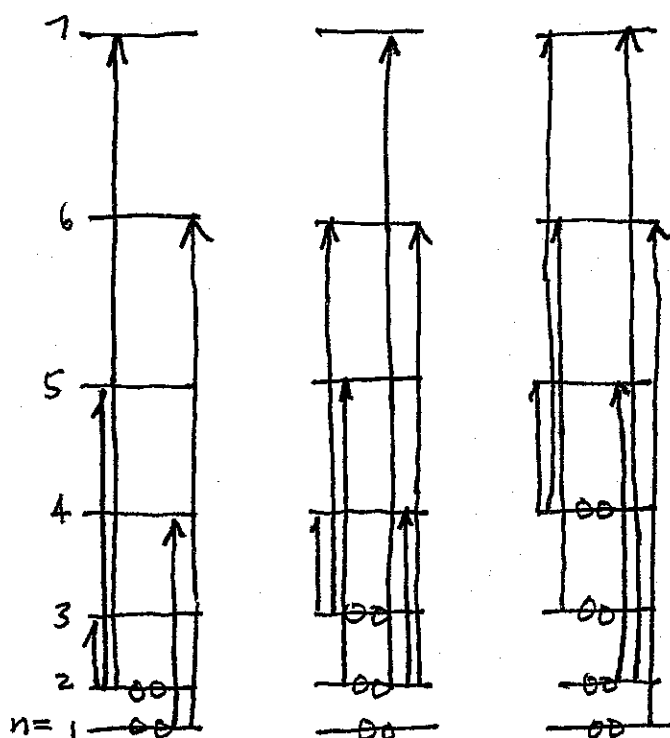
$$= 0 + 0 + 0 + x_{23} \text{ because of ORTHONORMAL condition.}$$

$$= \frac{-8a}{\pi^2} \left(\frac{6}{25} \right) \text{ from page 1 of exam.}$$

$$\text{square of the integral} = \left(\frac{-8a}{\pi^2} \frac{6}{25} \right)^2$$

We see from the matrix $\mathbb{X}$ given on page 1 of the exam, that the only transitions allowed are only those in which

- ONLY one electron has a different n in the ground and excited states
- the n values in the ground and excited configurations have to be those that have non-zero matrix elements in the $\mathbb{X}$ matrix on page 1
e.g., 12, 14, 16, ... 21, 23, 25, 27, ...



$$\Delta n = \pm 1, \pm 3, \pm 5, \dots$$

butadiene $a = 7.00 \text{ \AA}$ $\Delta E = \frac{(3^2 - 2^2)h^2}{8ma^2}$ $\Delta E = \frac{(4^2 - 3^2)h^2}{8ma^2}$ $\Delta E = \frac{(5^2 - 4^2)h^2}{8ma^2}$

$$\lambda = \frac{hc}{\Delta E} = \frac{8mca^2}{h(3^2 - 2^2)} \quad \lambda = 330 \left(\frac{9.80^2}{7} \right) \quad \lambda = 330 \left(\frac{12.60^2}{9} \right)$$

$$= 330 \text{ \AA}^{-1} \frac{(7 \text{ \AA})^2}{5} \quad \lambda = 4528 \text{ \AA} \quad \lambda = 5821 \text{ \AA}$$

$$\lambda = 3234 \text{ \AA}$$

Oscillator strength $= 1.085 \times 10^{-11} \frac{\mathbb{X}^2}{\lambda} = 1.085 \times 10^{-11} \frac{\left(\frac{6}{25}\right)^2 \left(\frac{8}{\pi^2}\right)^2 \frac{7^2}{3234} \times 10^{-8}}{0.622} \dots$ etc.

Note that the length L of the polyene is proportional to the number of carbons, m_C if we let the ends of the line extend to $(1/2)$ the C-C bond length beyond the last carbon nuclei. That is, $L = m_C r_{CC}$. Note also that the double occupancy afforded by the Pauli exclusion principle means that the smallest energy transitions corresponds to “from” quantum numbers $m_C/2$ and “to” quantum numbers $1+m_C/2$. That is, the energy of the transition is given by $\Delta E = \{(1+m_C/2)^2 - (m_C/2)^2\} h^2/8m_e L^2 = \{(1+m_C/2)^2 - (m_C/2)^2\} h^2/8m_e (m_C r_{CC})^2 = hc/\lambda_{abs}$

where the longest wavelength λ_{abs} is given by

$\lambda_{abs} = (m_C)^2/(1+m_C) 8m_e r_{CC}^2/h$ So we expect the longest wavelength in the absorption spectrum of the polyene having m_C carbons to be proportional to $(m_C)^2/(1+m_C)$, i.e., $4^2/5$, $6^2/7$, $8^2/9$ from butadiene to hexatriene to octatriene. A plot of the longest wavelength versus m_C will not be exactly linear, but nearly so.

Calculation of current densities using gauge-including atomic orbitals

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The current density for hexabenzocoronene is displayed in a plane parallel to the molecular framework and 1 bohr above it.

