Answers to Problem Set 6

1. First the function has to be normalized:

$$\Psi = 0.4\psi_1 + 0.5\psi_2 + c_3\psi_3 + 0.3\psi_4$$

by setting $\int \Psi^* \Psi d\tau = 1$ That ψ_1 , ψ_2 , ψ_3 , ψ_4 are all the eigenfunctions of an operator **T**, means they are individually normalized and they are also orthogonal to each other since they correspond to <u>different</u> eigenvalues. Integrating leads to $\int \Psi^* \Psi d\tau = 1 = (0.4)^2 \int \psi_1 * \psi_1 d\tau + (0.5)^2 \int \psi_2 * \psi_2 d\tau + (c_3)^2 \int \psi_3 * \psi_3 d\tau + (0.3)^2 \int \psi_4 * \psi_4 d\tau + all other integrals <math>\int \psi_1 * \psi_2 d\tau$ etc. are zero. $c_3 * c_3 = 0.5$, $c_3 = 1/\sqrt{2}$

The fraction of the time the eigenvalue a_3 will be observed is $c_3*c_3 = 0.5$ Also a_1 , a_2 , a_4 will be found the following fractions of time: $(0.4)^2$, $(0.5)^2$, $(0.3)^2$ The expected everyon of T is given by Postulate 3:

The expected average of T is given by Postulate 3:

 $\int \Psi^* T \Psi d\tau =$

 $\int (0.4\psi_1 + 0.5\psi_2 + 1/\sqrt{2}\psi_3 + 0.3\psi_4)^* T (0.4\psi_1 + 0.5\psi_2 + 1/\sqrt{2}\psi_3 + 0.3\psi_4) d\tau$ $= (0.4)^2 a_1 + (0.5)^2 a_2 + 1/2 a_3 + (0.3)^2 a_4$

The expected average of T^3 is given by Postulate 3: $\int \Psi^* T^3 \Psi d\tau =$ $\int (0.4\psi_1 + 0.5\psi_2 + 1/\sqrt{2}\psi_3 + 0.3\psi_4)^* T^3 (0.4\psi_1 + 0.5\psi_2 + 1/\sqrt{2}\psi_3 + 0.3\psi_4) d\tau$ $= (0.4)^2 (a_1)^3 + (0.5)^2 (a_2)^3 + 1/2 (a_3)^3 + (0.3)^2 (a_4)^3$ because $T^3\psi_1 = T^2 a_1\psi_1 = a_1 T^2\psi_1 = (a_1)^2 T\psi_1 = (a_1)^3\psi_1$ and so on.

2. Particle on a line from x=0 to x=a is described by $\Psi_{t0}=Nx(a-x)$ at t_0 (a) The only possible outcomes of a measurement of energy are the eigenvalues of the Hamiltonian operator for a particle on a line system, which are $n^2h^2/8ma^2$, where n = 1,2,3,4, etc.

Which eigenvalues would be observed will depend on which eigenfunctions are included in $\Psi_{t0} = Nx(a-x)$.

(b) Average of single measurements from multiple copies is given by: Ave = $\int \Psi_{t0} * \mathcal{H} \Psi_{t0} dx / \int \Psi_{t0} * \Psi_{t0} dx = \int N * x(a-x) \mathcal{H} N x(a-x) dx / \int N^2 x^2 (a-x)^2 dx$ - $(\hbar^2/2m) d^2/dx^2 x(a-x) = ? d/dx [x(a-x)] = -x + (a-x) = a-2x; d/dx [a-2x] = -2$ - $(\hbar^2/2m) d^2/dx^2 x(a-x) = -(\hbar^2/2m) (-2)$ Ave = $-(\hbar^2/2m) \int -2x(a-x) dx / \int x^2 (a-x)^2 dx = ?$ $\int x(a-x) dx = |ax^2/2 - x^3/3|_0^a = a^3/2 - a^3/3 = a^3/6$ $\begin{aligned} \int x^2 (a-x)^2 dx = \int a^2 x^2 - 2ax^3 + x^4 dx &= \left| ax^3/3 - 2ax^4/4 + x^5/5 \right|_0{}^a = a^4/3 - a^5/2 + a^5/5 = a^5/30 \\ \text{Ave} = -(\hbar^2/2m) - 2[a^3/6] / a^5/30 = +5(\hbar^2/ma^2) \\ \text{Normalization of } \Psi_{t0}: \int N^2 x^2 (a-x)^2 dx = 1 = N^2 a^5/30 \text{ thus, } N^2 = 30/a^5 \\ \text{and } N = [30/a^5]^{\frac{N}{2}} \\ \text{(c) Expand in terms of the complete set of eigenfunctions of the system:} \\ \Psi_{t0} = \sum_n C_n \psi_n \\ \text{Operate on both sides with } \int \psi_i * dx, \text{ to get } C_i \\ \int \psi_n * [30/a^5]^{\frac{N}{2}} x(a-x) dx = C_n \\ \text{since all other integrals in the sum } \sum_i C_i \int \psi_n * \psi_i dx \text{ vanish for } i \neq n \text{ since the eigenfunctions of the system are orthonormal.} \\ \int \psi_n * [30/a^5]^{\frac{N}{2}} x(a-x) dx = C_n = [2/a]^{\frac{N}{2}} [30/a^5]^{\frac{N}{2}} \int \sin(n\pi x/a) x(a-x) dx \\ C_n = [60/a^6]^{\frac{N}{2}} \left\{ \int ax \sin(n\pi x/a) dx - \int x^2 \sin(n\pi x/a) dx \right\} \end{aligned}$

Use known integrals

 $\int x \sin(bx) dx = (1/b^2) \sin(bx) - (x/b) \cos(bx)$ and $\int x^2 \sin(bx) dx = -[(b^2x^2 - 2)/b^3] \cos(bx) + 2x \sin(bx)/b^2$

After the smoke clears in the integrations, we find that when n is odd, we get $C_{odd} = 0$ when n is even, we get

 $C_n = [60/a^6]^{\frac{1}{2}} 4a^3/(n\pi)^3$

3. Consider a particle of mass M constrained to move on a circle of radius R where its potential energy is zero. The particle is in a physical state that is described by $F(\phi) = A\{\cos 2\phi + 2\cos 3\phi\}$. Determine the results of the following sets of experiments on this system, that is, determine the typical outcomes of the experiments, the average values of the results:

(a) The z component of the angular momentum of the system is measured

| Derivation of predictions here: | Observed |
|---|--------------|
| $F(\phi)$ is not one of the eigenfunctions of L_z or \mathcal{H} . | values here: |
| Expand $F(\phi)$ in terms of the complete orthonormal set of functions | 2 ħ |
| $(1/\sqrt{2\pi})\exp[ik\phi]$. $F(\phi) = \sum_k c_k \Psi_k(\phi)$ and find the coefficients: | -2 ħ |
| $c_{k} = \int_{0}^{2\pi} \Psi_{k}^{*}(\phi) F(\phi) d\phi = \int_{0}^{2\pi} \Psi_{k}^{*}(\phi) A\{\cos 2\phi + 2\cos 3\phi\} d\phi$ | 3 ħ |
| Since we can write $\cos 2\phi = \frac{1}{2} [\exp(i2\phi) + \exp(-i2\phi)]$ and | 3 ħ |
| $\cos 3\phi = \frac{1}{2} [\exp(i\beta\phi) + \exp(-i\beta\phi)], \text{ then}$ | 3 ħ |
| $c_{k} = \int_{0}^{2\pi} \Psi_{k}^{*}(\phi) \frac{1}{2} A \{ \exp(i2\phi) + \exp(-i2\phi) + 2\exp(i3\phi) + 2\exp(-i3\phi) \} d\phi$ | 3 ħ |
| $c_{k} = (1/\sqrt{2\pi})^{-1} \sqrt{2\pi} A \int_{0}^{2\pi} \Psi_{k}^{*}(\phi) \{\Psi_{2}(\phi) + \Psi_{-2}(\phi) + 2\Psi_{3}(\phi) + 2\Psi_{-3}(\phi)\} d\phi$ | -3 ħ |

| $c_k = (1/\sqrt{2\pi})^{-1} \sqrt{2} A \{ \delta_{k,2} + \delta_{k,-2} + 2\delta_{k,3} + 2\delta_{k,-3} \} $ in shorthand | -3 ħ |
|---|--------------------|
| where $\delta_{k,2}=1$ if k=2 or else it is zero, since it was an orthonormal set | -3 ħ |
| A can be obtained by integration $\int_0^{2\pi} F^*(\phi)F(\phi) d\phi$: $A = [5\pi]^{-1/2}$ | -3 ħ |
| $c_{k} = (10)^{-1/2} \{ \delta_{k,2} + \delta_{k,-2} + 2\delta_{k,3} + 2\delta_{k,-3} \}$ | Average = $0\hbar$ |
| $c_2^2 = c_{-2}^2 = (1/10)$ $c_3^2 = c_{-3}^2 = (4/10)$ | _ |

Observed values here:

 $2^{2} (\hbar^{2}/2MR^{2})$

 $(-2)^2 (\hbar^2 / 2MR^2)$

 $\frac{3^2}{3^2} (\hbar^2 / 2MR^2)$ $3^2 (\hbar^2 / 2MR^2)$

 $3^2 (\overline{\hbar^2/2MR^2})$

 $\frac{3^{2} (\hbar^{2}/2MR^{2})}{(-3)^{2} (\hbar^{2}/2MR^{2})}$ $\frac{(-3)^{2} (\hbar^{2}/2MR^{2})}{(-3)^{2} (\hbar^{2}/2MR^{2})}$ $\frac{(-3)^{2} (\hbar^{2}/2MR^{2})}{(-3)^{2} (\hbar^{2}/2MR^{2})}$

Average = $8 (\hbar^2/2MR^2)$

(b) The energy of the system is measured

Derivation of predictions here:

The average value is obtained by $\int_0^{2\pi} F^*(\phi) (L_z \text{ or } \mathcal{H}) F(\phi) d\phi$: which by above algebra leads to $\sum_k c_k^2 \bullet k\hbar$ or $\sum_k c_k^2 \bullet k^2(\hbar^2/2MR^2)$ The observed values should be 10% of the time the eigenvalue for k=2, 10% of the time the eigenvalue for k=-2, 40% of the time the eigenvalue for k=3, 40% of the time the eigenvalue for k=-3, according to the probabilities given by the corresponding c_k^2

4. A particle of mass *m* in a potential well (with infinitely high walls) in the x dimension is known to be in either the n = 2 or n = 3 eigenstates with equal probability. The eigenfunctions of these states are $\psi_2(x) = (2/a)^{1/2} \sin [2\pi x/a]$ and $\psi_3(x) = (2/a)^{1/2} \sin [3\pi x/a]$, respectively.

(a) Write an appropriate wavefunction Ψ for the system that reflects our knowledge of the state of the system.

According to the conditions of the problem both states are equally probable, thus we need to have the wavefunction be a superposition of $\psi_2(x)$ and $\psi_3(x)$ with coefficients whose absolute squares are equal. $\Psi(x) = c_2\psi_2(x) + c_3\psi_3(x)$ such that $c_2^2 = c_3^2 = 1/2$ since $c_2^2 + c_3^2 = 1$ (normalization). Therefore, $\Psi(x) = (1/\sqrt{2})\{(2/a)^{1/2} \sin [2\pi x/a] + (2/a)^{1/2} \sin [3\pi x/a]\}$ or $\Psi(x) = (1/\sqrt{2})\{(2/a)^{1/2} \sin [2\pi x/a] - (2/a)^{1/2} \sin [3\pi x/a]\}$

(b) *What energies* might be obtained if the energy of the particle is measured? The energy eigenvalues $2^{2}h^{2}/8ma^{2}$ and $3^{2}h^{2}/8ma^{2}$ only.

(c) *Determine the expected average* of a series of measurements of the energy of the particle.

Postulate 3 says the expected average is $\langle E \rangle = \int_0^a \Psi^*(x) \mathcal{H} \Psi(x) dx$, since we have already normalized $\Psi(x)$.

 $\begin{aligned} \langle \mathbf{E} \rangle &= \int_0^a (1/\sqrt{2}) \{ \psi_2(\mathbf{x}) + \psi_3(\mathbf{x}) \}^* \, \mathcal{H}(1/\sqrt{2}) \{ \psi_2(\mathbf{x}) + \psi_3(\mathbf{x}) \} d\mathbf{x} \\ &= (1/2) \int_0^a \{ \psi_2^* \, \mathcal{H} \psi_2 d\mathbf{x} + \psi_2^* \, \mathcal{H} \psi_3 d\mathbf{x} + \psi_3^* \, \mathcal{H} \psi_2 d\mathbf{x} + \psi_3^* \, \mathcal{H} \psi_3 d\mathbf{x} \} \\ &= (1/2) \{ \mathbf{E}_2 \int_0^a \psi_2^* \psi_2 d\mathbf{x} + \mathbf{E}_3 \int_0^a \psi_2^* \psi_3 d\mathbf{x} + \mathbf{E}_2 \int_0^a \psi_3^* \psi_2 d\mathbf{x} + \mathbf{E}_3 \int_0^a \psi_3^* \psi_3 d\mathbf{x} \} \\ \langle \mathbf{E} \rangle &= (1/2) \{ \mathbf{E}_2 + \mathbf{E}_3 \} = (1/2) \{ 2^2 + 3^2 \} h^2 / 8 ma^2 \end{aligned}$

For the - combination the results are the same as above: $\langle E \rangle = \int_0^a (1/\sqrt{2}) \{ \psi_2(x) - \psi_3(x) \}^* \mathcal{H}(1/\sqrt{2}) \{ \psi_2(x) - \psi_3(x) \} dx$ $= (1/2) \{ E_2 \int_0^a \psi_2^* \psi_2 dx - E_3 \int_0^a \psi_2^* \psi_3 dx - E_2 \int_0^a \psi_3^* \psi_2 dx + E_3 \int_0^a \psi_3^* \psi_3 dx \}$ $\langle E \rangle = (1/2) \{ E_2 + E_3 \} = (1/2) \{ 2^2 + 3^2 \} h^2 / 8ma^2$

(d) *Write the equation* that shows how the <u>expected mean square deviation</u> of any series of measurements of the energy of the particle can be calculated.

The operator for the square of the deviation in measurements of energy is Op = $(\mathcal{H} - \langle E \rangle)^2$ Postulate 3 gives the average, thus the mean square deviation = $\int_0^a (1/\sqrt{2}) \{\psi_2(x) + \psi_3(x)\}^* (\mathcal{H} - \langle E \rangle)^2 (1/\sqrt{2}) \{\psi_2(x) + \psi_3(x)\} dx$

(e) *Carry out the solution* of (e), and then from the final result, determine the expected standard deviation of the series of measurements.

mean square dev = $\int_0^a (1/\sqrt{2}) \{\psi_2(x) + \psi_3(x)\}^* (\mathcal{H} - \langle E \rangle)^2 (1/\sqrt{2}) \{\psi_2(x) + \psi_3(x)\} dx$ = $\int_0^a (1/\sqrt{2}) \{\psi_2(x) + \psi_3(x)\}^* (\mathcal{H}^2 - 2\langle E \rangle \mathcal{H} + \langle E \rangle^2) (1/\sqrt{2}) \{\psi_2(x) + \psi_3(x)\} dx$ Note that $\mathcal{H}^2 \psi_2(x) = \mathcal{H} E_2 \psi_2(x) = E_2 \mathcal{H} \psi_2(x) = E_2^2 \psi_2(x)$ and $\langle E \rangle \mathcal{H} \psi_2(x) = \langle E \rangle E_2 \psi_2(x)$ and $\langle E \rangle^2 \psi_2(x) = \langle E \rangle^2 \psi_2(x)$ since $\langle E \rangle^2$ is a number. mean square dev = $(1/2) \int_0^a \{\psi_2 + \psi_3\}^* (E_2^2 \psi_2 + E_3^2 \psi_3 - 2\langle E \rangle E_2 \psi_2 - 2\langle E \rangle E_3 \psi_3 + \langle E \rangle^2 \psi_2 + \langle E \rangle^2 \psi_3 \} dx$ = $(1/2) \{E_2^2 + E_3^2 - 2\langle E \rangle (E_2 + E_3) + 2\langle E \rangle^2 \} = (1/2) \{E_2^2 + E_3^2 - 4\langle E \rangle^2 + 2\langle E \rangle^2 \}$ = $(1/2) \{E_2^2 + E_3^2\} - \langle E \rangle^2 = [(1/2) \{2^4 + 3^4\} - \{(2^2 + 3^2)/2\}^2] \{h^2/8ma^2\}^2$ = $[25/4] \{h^2/8ma^2\}^2$ standard deviation is the square root of this = $(5/2)(h^2/8ma^2)$ (e) Illustrate a typical table of results from 10 such measurements. *Fill in* the column "<u>Results</u>". What is the probability of each outcome?

| | Result | Deviation | • · · · · · · · · · · · · · · · · · · · | Probability |
|-----|--------------|-------------|---|-------------|
| 1 | 22h2/8ma2 | -2,5h2(Rma2 | - | 1/2 |
| 2 | 22h2/8ma2 | | ····· | |
| 3 | 3 h2/8ma2 | +2,5h2/8ma2 | | 1/2 |
| 4 | 22h2/ema2 | | | |
| 5 | 32 h2 (8ma2 | | · · · | |
| 6 | 32h2/8mg2 | | | |
| 7 | 3 h2/8ma2 | | | |
| 8 | 22 h2/2 maz. | | 7 | |
| 9 ' | 32 h2/8 ma2 | | | |
| 10 | 22h2/8ma2 | | | |
| Ave | 6.5 h2/8 mar | | | |



(h) Suppose an electron is contained in a two-dimensional potential well (with infinitely high walls) whose shape is that of a rectangular sheet with dimensions $a \times b$. Write the Schrodinger equation that needs to be solved for this system.

$$\mathcal{H} = (x,y) = -\frac{\hbar^2}{2m_e} \left(\frac{d^2}{dx^2} + \frac{d^2}{dy^2} \right) = E = E(x,y)$$

(i) Show that the method of separation of variables may be used to solve this problem, i.e., to find the eigenfunctions and eigenvalues.

 $\underline{t} \Psi(x, y) = F(x) \cdot G(y)$ $\frac{d^2}{2me} \left[\frac{d^2}{dx^2} F(x) \cdot G(y) + \frac{d^2}{dy^2} F(x) \cdot G(y) \right] = E F(x) \cdot G(y)$ Dividing both sides by F(x). G(y) after quarting $\frac{-\frac{t^2}{2me}\left[G(y),\frac{d^2F(x)}{dx^2} + F(x)\frac{d^2G(y)}{dy^2}\right]}{F(x)\cdot G(y)} = \frac{EF(x)\cdot G(y)}{F(x)\cdot G(y)} = E$ $-\frac{t^2}{2m_e} \left[\frac{d^2 F(x)}{dx^2} \cdot \frac{1}{F(x)} + \frac{1}{G(y)} \frac{d^2 G(y)}{dy^2} \right] = E$ Since each of these daw independently take any x values unrelated to g values, Atagyt the sum has always to be the same constant value E then each term in the sum knust itself be a constant, independent of the value of X or y. -th² d² F(x) = A -th² - 1 d² G(y) = B, such that A+B = t can solve this

(j) Given the results of your proof above, write down the possible energy eigenfunctions for an electron confined to a sheet with dimensions $a \times b$. Given the results of your proof above, write down the corresponding energy eigenvalues opposite the eigenfunction

 $\frac{-t^2}{2m_e F(x)} \frac{d^2 F(x)}{dx^2} = A \quad or \quad \frac{-t^2}{2m_e} \frac{d^2 F(x)}{dx^2} = A F(x)$ Solving $2me f(x) dx^{2} \qquad 2me dx^{2}$ eigenvalues are already known, as seen in part (b) g this exam : $A = n_{x} h^{2}/8ma^{2}$ $B = \frac{n_{x}^{2}h^{2}}{8mb^{2}}$ Both equations have the same torm, therefore the solutions are the same torm, therefore $E = A + B = \frac{n_{x}^{2}h^{2}}{8mb^{2}} + \frac{n_{y}^{2}h^{2}}{8mb^{2}}$ where $n_{x} = 1, 2, 3, \cdots$ Functions are therefore $F(x) \cdot G(y)$ or $I(x, y) = \sqrt{2} \sin(\frac{n_{x}T}{a} x) \cdot \sqrt{2} \sin(\frac{n_{y}T}{b} y)$