

Electron Coupling of Nuclear Spins. VI. Relative Signs of $J_{\text{gem}}^{\text{HH}}$, J_o^{HH} , and J_t^{HH} in (2.2)Metacyclophane*

H. S. GUTOWSKY AND CYNTHIA JUAN

Noyes Chemical Laboratory, University of Illinois, Urbana, Illinois

(Received March 22, 1962)

The high-resolution proton magnetic resonance spectrum of the methylene groups in (2.2) metacyclophane has been observed at 60 and 15 Mc. In this compound, the C—CH₂—CH₂—C groups are fixed in position with the dihedral angle between alkyl C—C—C bonds slightly less than the symmetrical, staggered 60° and with little distortion from tetrahedral of the other bond angles in the group. Therefore, the HCCH coupling constants should be characteristic of ethanic *trans* and *gauche* orientations and the CH₂ should be comparable with the *geminal* coupling in methane.

A complete analysis of the A₂X₂ and A₂B₂ type spectra, in combination with previous, unambiguous experimental findings that $|J_{\text{gem}}^{\text{HH}}| \cong |J_t^{\text{HH}}| > |J_o^{\text{HH}}|$, leads to the following assignments in (2.2) metacyclophane: $J_t^{\text{HH}} = \pm 12.3$, J_o^{HH} (the coupling of the "central" pair of *gauche* protons) = ± 3.2 , $J_{\text{gem}}^{\text{HH}} = \mp 12.0$, and $J_o^{\text{HH}} = \pm 4.0$, all ± 0.1 cps. The magnitudes of these constants agree well with previous valence-bond calculations for CH₄ and the ethanic HCCH group. However, the opposite signs found for the large *trans* and *geminal* constants disagrees with the theoretical prediction that both are positive. The observed difference in sign can not be attributed to substituent effects, angular distortions, or to motional averaging. It is concluded that one of the two sets of calculations is in error; the implications of this result are discussed.

I. INTRODUCTION

THIS paper is concerned with the comparison between experiment and valence-bond calculations of the σ -electron contribution to the proton-proton coupling constants of CH₂ and HCCH groups in high-resolution NMR spectra.¹ Such calculations predicted the *geminal* coupling $J_{\text{gem}}^{\text{HH}}$ to be +12.5 cps in methane,² and subsequent, more approximate, calculations³ for the coupling of *vicinal* protons in the HCCH ethanic fragment with tetrahedral HCC bond angles gave the *trans* coupling J_t^{HH} to be about +9.2 cps and the *gauche* J_o^{HH} , +1.7 cps. The magnitudes of these values agree well with experiment except that the *trans vicinal* constants observed for ethanic groups⁴ (and also both the *cis* and *trans* constants for ethylene³) are often about 50% larger than predicted.

In most cases, only the magnitudes of the coupling constants have been obtained from experiment; but increasing attention is being given to the importance of determining their relative signs. Several instances of substituted ethylenes have been reported^{5,6} in which the sign of $J_{\text{gem}}^{\text{HH}}$ (1 to 3 cps) is opposite to, and also

the same as, that of $J_{\text{cis}}^{\text{HH}}$ (5 to 11 cps) and $J_{\text{trans}}^{\text{HH}}$ (12 to 18 cps). These results are compatible with the valence-bond calculations for the CH₂ fragment,⁵ which neglect substituent and π -electron effects, and which predict that $J_{\text{gem}}^{\text{HH}}$ should become negative for HCH bond angles larger than about 120°. In a similar vein, for the *vicinal* coupling, experiment has yielded instances of substituted ethanes in which the sign of J_o^{HH} (1 to 3 cps) is opposite to, and also the same as, that of J_t^{HH} (10 to 16.5 cps).^{4,6} These results also are compatible with the predicted dependence of the coupling upon the dihedral angle ϕ ,

$$J_{\text{vic}}^{\text{HH}} \cong 9 \cos^2 \phi - 0.3, \quad (1)$$

in view of the neglect in the valence-bond calculations³ of substituent effects, and the assumption of tetrahedral HCC angles.

However, while the present work was in progress a question less easy to dispose of was raised by reports^{8,9} of relative sign determinations which conflict with the predictions^{2,3} that large values of $J_{\text{gem}}^{\text{HH}}$ and $J_{\text{vic}}^{\text{HH}}$ should both be positive. In diethyl sulfite, the two protons in the CH₂ groups are structurally nonequivalent because of the molecular asymmetry, so the spectrum is of the *abc*₃ type.¹ Because of this asymmetry, $J_{\text{gem}}^{\text{HH}}$ could be determined in the CH₂ group and its sign found with respect to $J_{\text{vic}}^{\text{HH}}$ averaged over the CH₃ group reorientations.⁸ The results are $J_{\text{gem}}^{\text{HH}} = \pm 10.45$ cps and $J_{\text{ac}}^{\text{HH}} = J_{\text{bc}}^{\text{HH}} = \mp 7.12$ cps. Similar results were found⁹ for the —CH₂CHX— group of several 1,3-dioxo, 5-X, cyclopentanes, in which the

¹ H. S. Gutowsky, V. D. Mochel, and B. G. Somers, J. Chem. Phys. **36**, 1153 (1962).

² F. Kaplan and J. D. Roberts, J. Am. Chem. Soc. **83**, 4666 (1961).

³ R. R. Fraser, R. U. Lemieux, and J. D. Stevens, J. Am. Chem. Soc. **83**, 3901 (1961). See also the results on epichlorohydrin by C. A. Reilly and J. D. Swalen, J. Chem. Phys. **35**, 1522 (1961).

* Acknowledgment is made to donors of The Petroleum Research Fund, administered by the American Chemical Society, for partial support of this research. The work also was supported by the Office of Naval Research.

¹ J. A. Pople, W. G. Schneider, and H. J. Bernstein, *High-Resolution Nuclear Magnetic Resonance* (McGraw-Hill Book Company, Inc., New York, 1959).

² M. Karplus, D. H. Anderson, T. C. Farrar, and H. S. Gutowsky, J. Chem. Phys. **27**, 597 (1957); M. Karplus and D. H. Anderson, *ibid.* **30**, 6 (1959).

³ M. Karplus, J. Chem. Phys. **30**, 11 (1959).

⁴ H. S. Gutowsky, G. G. Belford, and P. E. McMahon, J. Chem. Phys. **36**, 3353 (1962), and prior work cited there for substituted ethanes.

⁵ H. S. Gutowsky, M. Karplus, and D. M. Grant, J. Chem. Phys. **31**, 1278 (1959).

⁶ C. N. Banwell, N. Sheppard, and J. J. Turner, Spectrochim. Acta **16**, 794 (1960).

two slightly different *vicinal* constants are about 7 cps and opposite in sign to the 8 cps value for $J_{\text{gem}}^{\text{HH}}$. While these results indicate that something probably is intrinsically wrong with either the methane calculation² and/or that for the ethanic fragment,³ the compounds studied are such that substituent effects, angular distortions and motional averaging are important, and their neglect in the theoretical treatment might be responsible for the apparent discrepancies in the relative signs.

A more convincing example is provided by the proton spectrum of the $-\text{CH}_2\text{CH}_2-$ groups in (2.2) metacyclophane, the structure of which is sketched in Fig. 1. The spectrum of this compound was first reported by Wilson, Boekelheide, and Griffin,¹⁰ in a study of chemical shifts produced by ring currents, and for which only a partial analysis was made. They gave the spatial relationships of the protons on the basis of an assumed model which led to the methylene groups being twisted about the CH_2-CH_2 bond by about 18° from the symmetrically staggered conformation. However, an x-ray structural determination has been made for the solid compound,¹¹ and the C—C—C bond angles of the alkyl groups found to be $109^\circ 49'$ and $110^\circ 48'$. Moreover, from the carbon positions given we have calculated the dihedral angle between the alkyl C—C—C bonds and found it to be $59 \pm 1^\circ$, which means that the methylene protons are in virtually a symmetrical staggered conformation. Therefore, this compound has the advantage that angular distortion effects should be minimal for the $-\text{CH}_2\text{CH}_2-$ groups. Moreover, the latter are locked into position so that individual values can be obtained for J_t^{HH} and J_g^{HH} instead of some uncertain kind of average; and, finally, only carbon and hydrogen atoms are bonded to the methylene groups, so substituent effects should be small. For these reasons, a complete analysis of the $-\text{CH}_2\text{CH}_2-$ spectrum was undertaken.

II. EXPERIMENTAL RESULTS

The sample of (2.2)metacyclophane used in these experiments was very kindly provided by Wilson, Boekelheide, and Griffin.¹⁰ The high-resolution proton spectra were observed at room temperature using 10% solutions in CCl_4 . The spectra at 60 Mc/sec were run locally on Varian Associates HR-60 and A-60 spectrometers. The 15.083 Mc/sec spectrum was obtained through the courtesy of J. N. Shoolery at Varian Associates, where it was run on a V-4300 spectrometer system with a 12-in. magnet. The spectra were calibrated by the usual audiofrequency sideband method.¹

The over-all procedure used to determine the magnitudes and relative signs of the coupling constants in the $-\text{CH}_2\text{CH}_2-$ group is the following. At a resonance frequency of 60 Mc/sec, the chemical shift $\nu_0\delta$ between

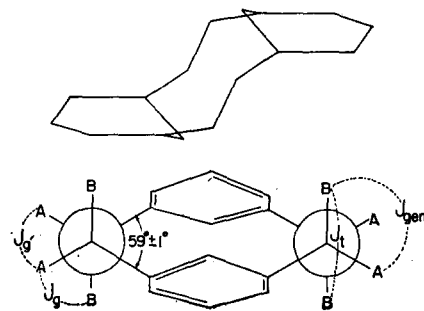


FIG. 1. The molecular geometry of (2.2)metacyclophane. The bottom part of the figure is a projection in the plane perpendicular to the C—C bond of the $-\text{CH}_2\text{CH}_2-$ groups; it includes definitions of the symbols used for the coupling constants.

the A_2 and B_2 sets of protons, defined in Fig. 1, is sufficiently large that the quite simple observed spectrum is a good approximation to the A_2X_2 type.¹ From it, the magnitudes of $\nu_0\delta$ and of the four coupling constants are determined readily, as well as the relative signs for each of two pairs of coupling constants. In part, the 60 Mc spectrum is easy to analyze because it is insensitive to one of the relative signs. However, the latter becomes important at lower resonance frequencies, where the spectrum is of the A_2B_2 type.¹ Therefore, the magnitudes and signs obtained from the 60 Mc spectrum were used to calculate 15.083 Mc spectra for the remaining relative sign permutations, and comparison of these with the observed spectrum completes the analysis.

We will not repeat here the theoretical expressions for the transition energies and intensities of A_2X_2 and A_2B_2 spectra, which are available in several publications¹ along with some methods for analyzing the spectra. However, determining the relative signs and magnitudes of all four coupling constants in such a four-spin system is sufficiently complex that it seems desirable to outline all of the steps followed in the present case.¹²

It is convenient to use parameters in the analysis which are the sum and difference of pairs of the coupling constants.¹ For our case, these parameters are

$$\begin{aligned} K &= \pm |J_t + J_{g'}| & N &= |J_g + J_{\text{gem}}| \\ M &= |J_t - J_{g'}| & L &= |J_g - J_{\text{gem}}|, \end{aligned} \quad (2)$$

where the coupling constants are defined in Fig. 1. In principle, each of the constants, and also the chemical shift, have both magnitude and sign. However, the spectrum is independent of the sign of the chemical shift and insensitive to inverting the signs of all coupling constants. This leaves the appearance of the spectrum governed by the magnitudes of the five

¹² D. M. Grant, R. C. Hirst, and H. S. Gutowsky, *J. Chem. Phys.* (to be published). This reference describes in considerable detail the nature and analysis of A_2B_2 and A_2X_2 spectra in general, and serves as a basis for the approach used here on (2.2) metacyclophane.

¹⁰ D. J. Wilson, V. Boekelheide, and R. Griffin, *J. Am. Chem. Soc.* **82**, 6302 (1960).

¹¹ C. J. Brown, *J. Chem. Soc.* **1953**, 3278.

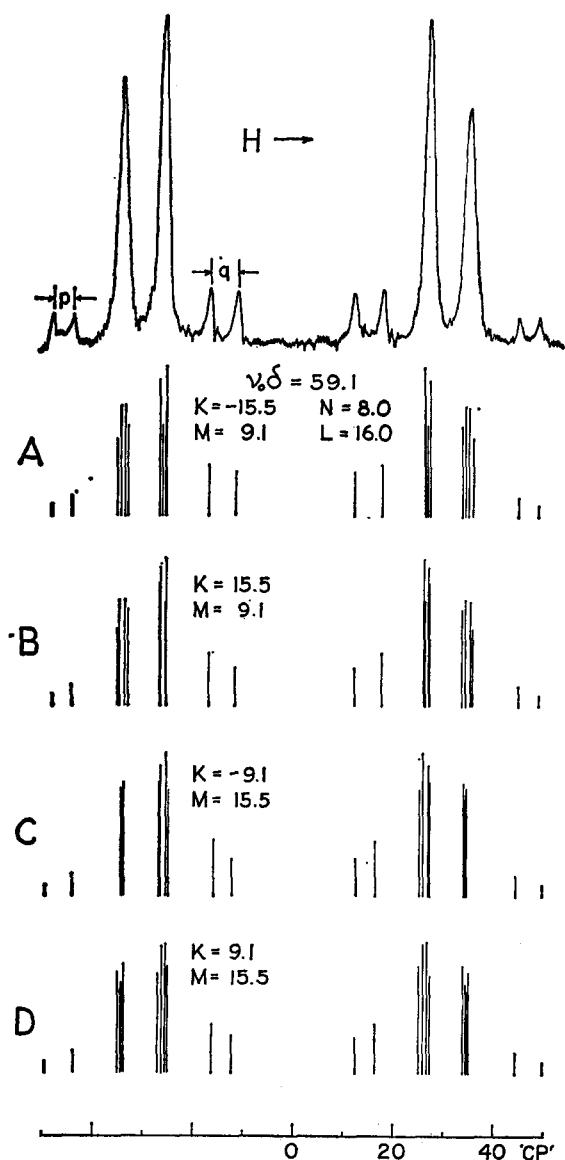


FIG. 2. At the top is the spectrum observed at 60 Mc for the $-\text{CH}_2\text{CH}_2-$ group protons in (2,2)metacyclophane. The line spectra shown below were calculated for the known values of N and L , 8.0 and 16.0 cps, and the four possible sign permutations $K = \pm 15.5$, $M = 9.1$, and $K = \pm 9.1$, $M = 15.5$ cps. Only the first two, A and B, fit the splittings p and q .

quantities and by the three relative signs of the four coupling constants.

In terms of the parameters K , L , M , and N , which we treat as positive quantities except for K in the one circumstance noted below, the relative signs of each pair of coupling constants in Eq. (2) are determined by the relative values of the corresponding two parameters. Thus, if J_g and J_{gem} have the same sign, $N > L$; and if opposite, $N < L$. Identical relations involving K and M hold for J_t and $J_{g'}$. In addition, the spectrum is sensitive to the actual relative signs of $K = (J_t + J_{g'})$ and $N = (J_g + J_{gem})$. Whether the observed spectrum is fitted by K positive or negative, while treating N as

positive, determines the third relative sign. If K negative applies, then the constant of largest magnitude in K is of opposite sign to the constant of largest magnitude in N , while they are of the same sign for a positive K . Finally, the magnitudes of the coupling constants are obtained by means of Eq. (2) from the numerical (positive) values for K , L , M , and N ; however, the spectrum alone does not tell which constant is which within each pair and supplemental information about the relative magnitudes of the constants is required to complete the assignment.

A. 60 Mc Spectrum

The proton spectrum observed at 60 Mc is shown in Fig. 2. To a first approximation it is of the A_2X_2 type,¹³ consisting of "mirror image" A_2 and X_2 multiplets whose centers are separated by $[(\nu_0\delta)^2 + N^2]^{\frac{1}{2}}$. In general, each of these two multiplets has ten lines, two quartets and a doublet with a common center. The splitting of the outer pairs of lines in one quartet is K , and of the other, M ; while the splitting of the inner pair is $(K^2 + L^2)^{\frac{1}{2}} - K$ and $(M^2 + L^2)^{\frac{1}{2}} - M$, respectively, for the two quartets. The lines of the doublet are the strongest transitions; their splitting is N .^{1,12}

In the observed spectrum, the A_2 and X_2 multiplets have two rather broad, very strong lines at the center, with two weaker lines at each side. Obviously, the inner lines of the two quartets fall on top of or very close to the strong N doublet, and only the outer lines of the quartets are resolved. Therefore, the ~ 8 cps splitting of the strong center pair of lines undoubtedly is N . Also, the outer splittings of the two quartets are ~ 9 and ~ 15 cps but at this point it is uncertain which is K and which is M . These values, in combination with the expressions for the central splitting of the two quartets and their observed values of ~ 8 cps, give an unambiguous value for L of 15 ± 2 cps. Also, the separation between the centers of the two multiplets is approximately the chemical shift, which gives $\nu_0\delta \cong 60.3$ cps.

Except for refinement of these numerical values, no further information could be obtained from the 60 Mc spectrum if it were truly of the A_2X_2 type.^{1,12} However, second-order effects are visible in the intensities of the A_2 and X_2 multiplets, which are peaked slightly towards one another. Also, the multiplets are asymmetric in that the separation between the outer lines from each quartet is less for the pairs of lines at the two ends of the spectrum than for the two pairs nearest the center, e.g., $p < q$ as in Fig. 2. Because of these deviations, efforts were made not only to refine the

¹³ The x-ray structure of the solid, reported in reference 11, indicates a 1° difference in the C—C bond angles at the two ends of the $-\text{CH}_2\text{CH}_2-$ groups. This could reduce the symmetry of the four-spin system from A_2X_2 to $ABXY$. However, the NMR spectra contain no evidence of such asymmetry, which may be too small for detection or else effectively absent in the liquid phase.

numerical values of the parameters but also to extend the analysis.

Two major aspects of the spectrum remain to be determined. First, which of the ~ 9 and ~ 15 cps values is K and which is M ? Secondly, does the observed spectrum correspond to K positive or negative? Therefore, spectra were calculated, and compared with experiment, for all four remaining possible relative sign permutations, $K \cong \pm 9$, $M \cong 15$; and $K \cong \pm 15$, $M \cong 9$. These calculations were made with the University of Illinois electronic digital computer, Illiac,¹⁴ using a program written for the general six-spin system.¹⁵ More refined parameters were obtained by varying them systematically, comparing the resulting calculated spectra with experiment, and then interpolating. In this manner, the following best-fit numerical values were obtained: $\nu_0\delta = 59.1$, $N = 8.0$, $L = 16.0$, and K or $M = 9.1$ or 15.5 , all in cycles per second. More important, the calculated spectra given in Fig. 2 for the four possible sign permutations show that although the spectrum observed at 60 Mc is too insensitive to the sign of K for its determination, the asymmetry in the splittings p and q in Fig. 2 is governed by the relative magnitudes of K and M . In order to have $p < q$ as observed, it is necessary to have $K > M$,¹⁶ which requires that K be 15.5 and M , 9.1 cps. Of the four calculated spectra in Fig. 2, A and B, with $K = \mp 15.5$ and $M = 9.1$ cps, agree with experiment, while C and D, with $K = \mp 9.1$ and $M = 15.5$ cps do not.

B. 15.083 Mc Spectrum

In order to determine the relative sign of K , the spectrum was observed at 15.083 Mc, where the chemical shift is small enough (14.86 cps) that the spectrum is definitely A_2B_2 . The results are shown in Fig. 3. A complete analysis might very well have been obtained from a single, A_2B_2 -type spectrum at say 30 Mc or even 40 Mc.¹² However, the analysis of a spectrum run at a high, as well as a low frequency minimizes the possibility of error. Figure 3 includes the observed spectrum and also spectra calculated for the two remaining sign permutations, $K = \pm 15.5$ cps. It is seen that there is excellent agreement between experiment and the spectrum calculated for $K = -15.5$ cps, and very poor agreement for $K = 15.5$ cps. As a check, spectra not shown were also calculated for $K = \pm 9.1$ and $M = 15.5$ cps, and they too do not agree with experiment.

Therefore, we conclude that the parameters which apply to the $-\text{CH}_2\text{CH}_2-$ group are

$$K = -15.5, \quad M = 9.1, \quad N = 8.0, \quad L = 16.0 \text{ cps.} \quad (3)$$

¹⁴ We are indebted to the staff of the Digital Computer Laboratory for their assistance with the calculations.

¹⁵ This program was written by G. G. Belford.

¹⁶ The relation found between the splitting inequality and the relative (absolute) values of K and M holds over the considerable range of parameters for which spectra were calculated. However, the relation is reversed if one takes $N > L$ instead of $N < L$ as it must be here; see, e.g., the calculated spectra in reference 12.

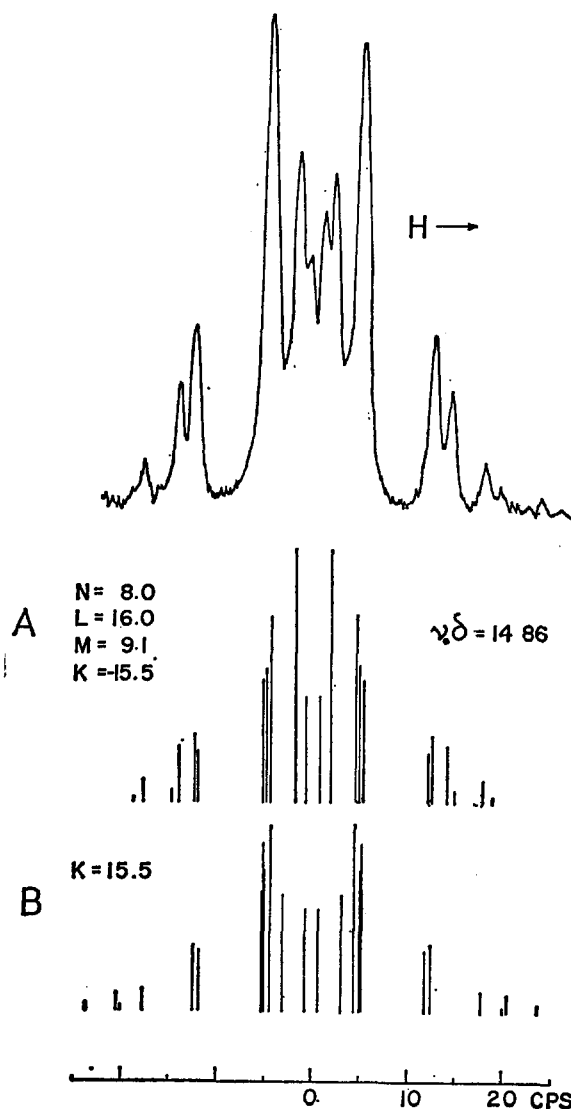


FIG. 3. At the top is the spectrum observed at 15.083 Mc for the $-\text{CH}_2\text{CH}_2-$ group protons in (2,2)metacyclophane. The line spectra shown below were calculated for the two remaining sign permutations, i.e., for K positive and negative. The bottom spectrum, for K positive disagrees greatly with the center portion and also the wings of the observed spectrum.

Upon combining these results with the definitions in Eq. (2) we find from N and L that J_θ and J_{gem} are 12.0 and 4.0 or 4.0 and 12.0 cps. Moreover, they are of opposite signs because $N < L$. From K and M , J_t and $J_{\theta'}$ are 12.3 and 3.2 or 3.2 and 12.3 cps. Also, they are of the same sign because $K > M$. (Here, both K and M must be treated as positive quantities.) Furthermore, since K and N actually have opposite signs, the largest constant of the K , M pair (12.3 cps) is of opposite sign to the largest constant of the N , L pair (12.0 cps). The assignment is completed by introducing the inequality $|J_t| > |J_{\theta'}|$, which is known with certainty from the NMR studies of substituted ethanes,^{4,6} and the inequality $|J_{\text{gem}}| > |J_\theta|$ which is equally certain

from the experimental results^{4,6} for substituted ethanes in combination with those on J_{gem} in methane² and in substituted methanes having HCH angles near tetrahedral.⁵ The final assignment is

$$\begin{aligned} J_t^{HH} &= \pm 12.3 \text{ cps}, & J_{\text{gem}}^{HH} &= \mp 12.0 \text{ cps}, \\ J_o^{HH} &= \pm 3.2, & J_q^{HH} &= \pm 4.0, \end{aligned} \quad (4)$$

where the probable errors of the numerical values are about ± 0.1 cps.

III. DISCUSSION

Our finding of large values of opposite sign for J_t^{HH} and J_{gem}^{HH} leaves little room for maneuver. The closeness of the 12.0 ± 0.1 cps value for J_{gem}^{HH} to the 12.4 ± 0.6 cps observed in methane² indicates that it is not affected by angular distortion and substituent effects. The small difference between the 3.2 and 4.0 cps values for J_o^{HH} and J_q^{HH} is compatible with the C—C—C dihedral angle being slightly less than the 60° for a symmetric, staggered C—CH₂CH₂—C group, which is suggested by the x-ray data for the solid.¹¹ Also, this would account for the value of 12.3 cps for J_t^{HH} being on the small side of values found for substituted ethanes.^{6,7} Therefore, our results, as well as the less conclusive earlier studies^{8,9} show that either the calculation on CH₄² or that on the ethanic (and probably also on the ethylenic) fragment³ is in error.

The implications of this finding are unfortunate. For example, comparisons with experiment of theoretical calculations of π -electron contributions to long-range proton-proton coupling constants in unsaturated straight chain compounds were based on the assumption that J_{vic}^{HH} is positive.¹⁷ This comparison, as well as many others, is now in doubt, and it is essential to clarify the problem. One can speculate as to which of the calculations is most likely to be in error, if not both. In *some* ways, the calculations for the HCCH fragment present the best opportunities for error. The calculations are more complex than for CH₄

(or CH₂), and it is possible for example that the non-neighboring-atom exchange integrals should not have been neglected.³

A more direct approach to the problem would be to determine J_{vic}^{HH} and/or J_{gem}^{HH} with respect to $J_{\text{C}^{13}\text{-H}}$. The calculation of the latter is considerably simpler than for J_{HH} , and there is little doubt but that J_{CH} is positive.^{18,19} At least such relative sign determinations would suggest which of the J_{HH} calculations should be redone first. Fortunately, it appears feasible to determine the relative signs of J_{CH} , J_{gem}^{HH} , and J_{vic}^{HH} by the sort of approach used here and also by double resonance methods.

In fact, analysis of the P³¹ and proton spectra observed for diphosphine H₂PPH₂ has given results,²⁰ bearing at least indirectly on our problem. For this compound, J_{gem}^{HH} and J_{PP} were found to have values of 108.2 ± 0.2 and 12 ± 4 cps, respectively, and to be of the same sign. Moreover, J_{vic}^{HH} (*cis* and *trans*) has values of 10.5 ± 0.2 and 6.8 ± 0.2 cps which are of the same sign but opposite to J_{gem}^{HH} and J_{PP} . By analogy to the results of the HCCH calculations,³ it was assumed²⁰ that J_{vic}^{HH} was positive in diphosphine, which, of course, made J_{PP} and J_{gem}^{HH} negative. A negative value for J_{PP} is surprising because the coupling between directly bonded atoms due to the usually dominant contact term is positive. In view of the present findings it *may be* somewhat more plausible to take J_{PP} as positive, which leads to J_{gem}^{HH} positive and J_{vic}^{HH} (*cis* and *trans*) negative. Moreover, a positive value of ~ 12 cps for J_{gem}^{HH} in diphosphine, and in phosphine,²⁰ is consistent with the small, 94° HPH bond angle and the valence-bond calculations of positive *geminal* coupling constants in CH₄² and in SiH₄.²¹ But the consistency does not prove that the latter are correct.

¹⁸ M. Karplus and D. M. Grant, Proc. Natl. Acad. Sci. U. S. 45, 1269 (1959).

¹⁹ H. S. Gutowsky and C. S. Juan, J. Am. Chem. Soc. 84, 307 (1962).

²⁰ R. M. Lynden-Bell, Trans. Faraday Soc. 57, 888 (1961).

²¹ J. C. Schug and H. S. Gutowsky, J. Chem. Phys. (to be published).

¹⁷ M. Karplus, J. Chem. Phys. 33, 1842 (1960).